

EPR study of ZnO:Co thin films grown by the PLD method

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Abstract. We have studied magnetic properties of zinc-oxide composite doped with a relatively high concentration (4%) of Co ions. Samples were obtained by the pulsed laser deposition (PLD) method with different parameters of the laser ablation. For various samples, the power of laser was changed as well as the temperature of sample heating. Electron paramagnetic resonance (EPR) measurements were carried out and temperature dependence of the EPR spectra was obtained. Analysis of the temperature dependences of the integral intensity of EPR spectra was carried out using the Curie-Weiss law. Fitting yields the following values $\theta_{CW} = 150$ K and $C = 1.67^{10}$, where C is the Curie constant, and θ_{CW} is the paramagnetic Curie temperature.

Key words: electron paramagnetic resonance (EPR) • diluted magnetic semiconductor (DMS) • ZnO:Co • Curie temperature

Introduction

Huge interest in ferromagnetic semiconductors with a wide energy gap is associated with the capabilities of their various applications such as in spintronics and the so-called semi-transparent electronics [4]. Due to the rare physical properties (the energy gap of about 3.4 eV at room temperature, the exciton binding energy of about 60 meV), ZnO is used for devices based on semiconductor structures [30, 31], such as varistors [34] or sensors [32]. Relatively simple production technology of ZnO crystals gives us the hope for lower cost of production equipment based on ZnO in comparison to the popular GaN. In addition, because of the ferromagnetic properties, this oxide can be used as a diluted magnetic semiconductor (DMS).

DMS are of interest for studying their unique spintronics properties with possible technological applications [8, 15], such as spin-spin exchange interaction between the localized magnetic moments and the band electrons [6]. This property of DMS has potential applications in the spin-dependent semiconductor electronics [2].

Some subsequent theoretical approach using density functional theory (DFT) [20, 21] and experimental [16, 26] works show that the n-type Co-doped ZnO also possesses room temperature ferromagnetism.

Some authors have reported ferromagnetism in ZnO doped with transition metal systems at the Curie temperatures T_C ranging from 30 to 550 K [3, 11, 14, 22,

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Received: 17 August 2012

Accepted: 4 January 2013

23, 26, 29] while others have found antiferromagnetic, spin glass, or paramagnetic behaviour [1, 12, 28].

The existence of ferromagnetic ordering in (Zn,Co)O has been theoretically proposed to be attributed to the double exchange interaction [20, 21] or to the RKKY interaction between Co ions [10]. Later calculation showed that the ground state of Co-doped ZnO is the spin glass without doping due to the short range interactions between transition metal atoms [13].

According to the theory of Dietl *et al.* [5], the p-d interactions are the reason for the long-range magnetic coupling. However, the investigated magnetic ZnO samples are either n-type conducting or insulating. The observed weak ferromagnetism in implanted ZnO, containing ferromagnetic nanocrystals and in PLD grown ZnO with diluted magnetic ions, is due to the ferromagnetic nanocrystals and the acceptor-like defects [33], respectively. For the diluted ZnCoO, it would be interesting to check if there is an interaction between the free electrons and d-electrons in the Co ions, and if this s-d interaction results in ferromagnetic coupling.

We report here the magnetic properties of zinc-oxide composite (ZnO) doped with Co. Electron paramag-

netic resonance (EPR) spectra have been measured and analyzed to extract information on the characteristics of the incorporation of the ions in the lattice.

Experimental

Samples of zinc-oxide doped with cobalt were obtained by the pulsed laser deposition (PLD) method with different parameters of the laser ablation. For various samples, the power of laser was changed (from 6 up to 600 impulses), as well as temperature of sample heating. Only some of the samples were heated at 573 K. The concentration of Co was found by the proportions of ZnO and CoO, $(\text{ZnO})_96$ and $(\text{CoO})_4$ to make a target of the PLD method (described in detail in work [25]). EPR measurements were performed with a modified X-band EPR spectrometer [24]. The temperature of the samples was controlled in the range of 100–370 K using a BRUKER liquid N gas flow cryostat.

Results and discussion

Transmission electron microscopy (TEM) images of the investigated samples of ZnO:Co are presented in Fig. 1. We included TEM image to show the morphology of the sample and present the material structure. The sample of ZnO:Co have porous structure with different graininess. The thickness of the layer, measured by scanning electron microscopy (SEM) was about 500 nm.

EPR spectra as a function of temperature for the annealed and non-annealed sample are shown in Figs. 2 and 3, respectively.

EPR spectra of annealed and non-annealed layers are similar with one difference, i.e. in the EPR spectrum of non-annealed layer appears an additional line with a value of $g = 1.96$ which can be associated with defects. Recently, it was concluded that the g -factor of

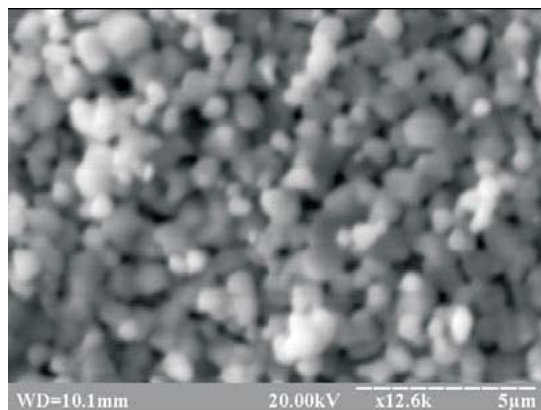


Fig. 1. Characteristics of the sample ZnO:Co, TEM image.

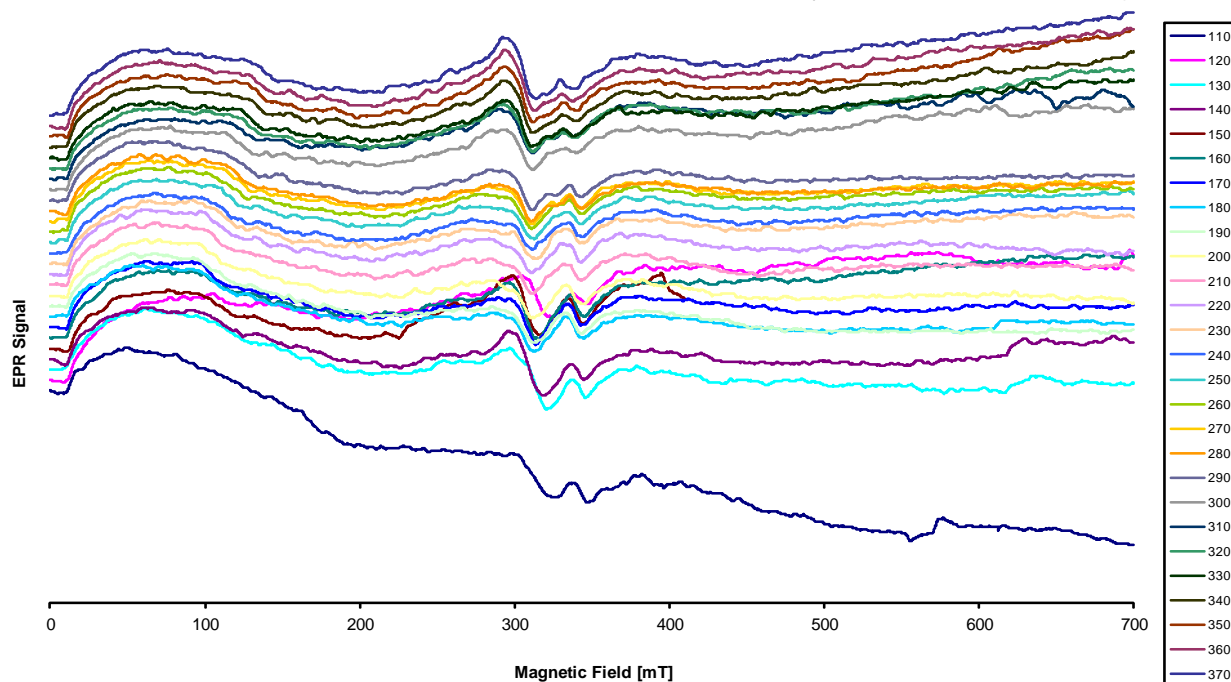


Fig. 2. EPR spectra of ZnCoO non-annealed sample (obtained at 600 pulses) at various temperatures.

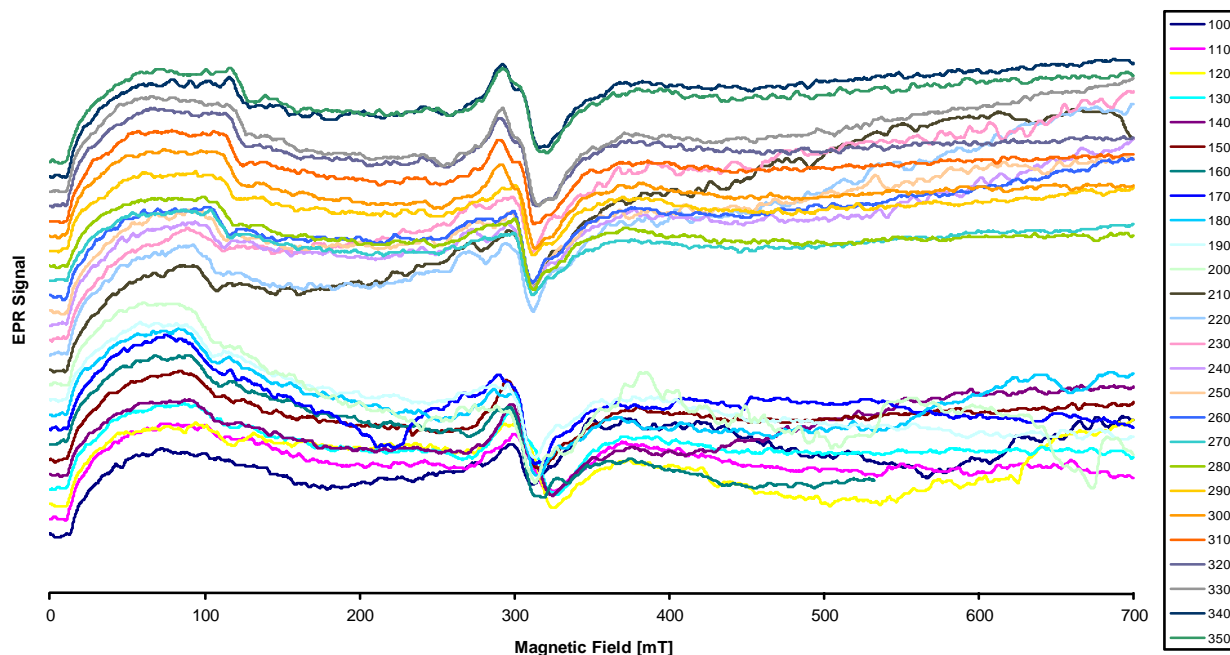


Fig. 3. EPR spectra of ZnCoO annealed sample (obtained at 600 pulses) at various temperatures.

1.96 is related to the oxygen vacancy, which was found by means of optically detected magnetic resonance [17, 27]. For samples obtained with a smaller number of pulses, the spectrum EPR are very similar, differing only in intensity. This is closely related to the thickness of the obtained layer, because for 600 pulses the greatest thickness was obtained, and, therefore, the signal is most intense. EPR spectra were analyzed for the layer of annealed and non-annealed ZnCoO.

The effective spectroscopic g -factor and the peak-to-peak line width (H_{pp}) of the resonance line were

Table 1. Numerical values of g -factor, H_r , H_{pp} and I

Temperature (K)	g -factor	H_r (mT)	H_{pp}	I
340	2.1950	302.06	18.46	246
330	2.2004	301.33	19.59	224
320	2.1968	301.82	20.32	225
310	2.2022	301.08	20.76	210
300	2.2055	300.63	23.88	197
290	2.2243	298.09	25.89	188
280	2.2294	297.41	26.75	177
270	2.2296	297.39	27.73	178
260	2.2511	294.54	31.96	165
250	2.2462	295.18	34.95	153
240	2.2538	294.18	34.56	142
230	2.2530	294.29	34.69	140
220	2.2609	293.26	34.18	160
210	2.2371	296.39	30.15	168
200	2.2328	296.96	27.81	177
190	2.2155	299.27	27.58	193
180	2.2163	299.16	26.66	200
170	2.1981	301.64	23.93	223
160	2.1647	306.29	19.23	265
150	2.1557	307.57	17.86	282
140	2.1590	307.10	22.39	304
130	2.1479	308.69	23.67	255
120	2.1074	314.62	21.53	208
110	2.1276	311.63	28.88	208

determined. Since the broad EPR line is asymmetric, the accuracy of parameters measured directly from the experimental spectrum is rather limited. Therefore, additionally the experimental line was fitted using the Lorentzian type curve, since such curves describe satisfactorily experimental EPR lines of DMS with manganese in high temperature range (see [18], and references therein). In this way we determined parameters for EPR lines such as the peak-to-peak line width (H_{pp}), the intensity (I) as well as the resonance field (H_r) (Table 1). Based on these data, we obtained the temperature dependencies of gyromagnetic factor $g(T)$, $H_r(T)$, and integral intensity as presented in Figs. 4 and 5, respectively. At fitting the Lorentz type curve, the lower field part of the EPR line was included because of its regular nature.

We used the Curie-Weiss law to analyze the temperature dependences of the integral intensity, which is directly proportional to the magnetic susceptibility χ . A linear increase of $\chi^{-1}(T)$ at higher temperatures can be fitted to the Curie-Weiss law [7, 9].

$$(1) \quad (\chi(T) - \chi_0)^{-1} = (T - \theta_{CW}) / C$$

where C is the Curie constant, θ_{CW} is the paramagnetic Curie temperature, and χ_0 is a temperature independent

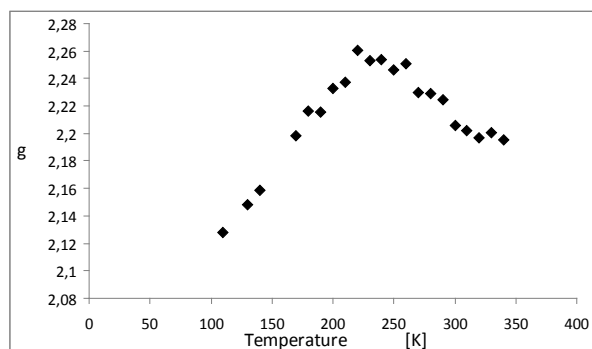


Fig. 4. Temperature dependence of g -factor for an annealed layer.

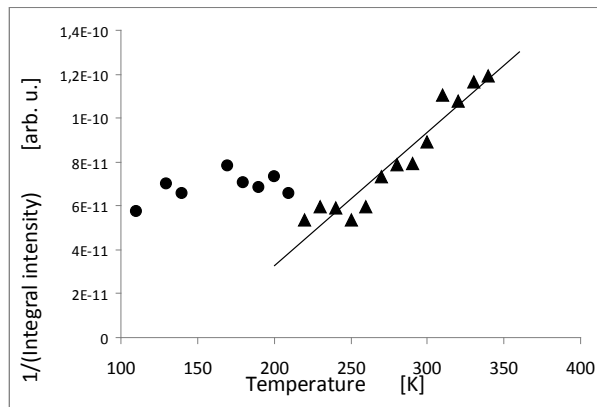


Fig. 5. Temperature dependence of the $1/(\text{integral intensity})$ for an annealed layer.

term to account for the diamagnetic host and any Pauli paramagnetism contribution. Figure 5 displays the temperature dependence of the quantity $(\chi(T) - \chi_0)^{-1}$. The line is a linear extrapolation illustrating the ferromagnetic (positive) Curie-Weiss temperature. Fitting yields the following values $\theta_{CW} = 150$ K and $C = 1.67^{10}$.

Conclusions

For annealed and non-annealed layers, the observed characteristic temperature around 220 K in which all EPR line parameters such as: peak-to-peak line width (H_{pp}), the intensity (I) as well as the gyromagnetic factor $g(T)$ are changed.

In summary, we have reported the X-band EPR studies of ZnO:Co. From the EPR lines we have determined the parameters: peak-to-peak line width (H_{pp}), intensity (I) as well as the resonance field (H_r). The results of temperature dependence of EPR spectra for the sample ZnO:Co and linear extrapolations to the Curie-Weiss law indicate the ferromagnetic interaction between the Co ions characterized by the Curie temperature 150 K.

By analyzing the EPR linewidth behaviour, the same as in [19], we claim that a combined effect of the exchange and dipolar broadening plays an important role in the mechanism of linewidth formation in this material.

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