# IMPROVED BOUNDS FOR SOLUTIONS OF $\phi$-LAPLACIANS 

Waldo Arriagada and Jorge Huentutripay<br>Communicated by Vicentiu D. Radulescu


#### Abstract

In this short paper we prove a parametric version of the Harnack inequality for $\phi$-Laplacian equations. In this sense, the estimates are optimal and represent an improvement of previous bounds for this kind of operators.


Keywords: Orlicz-Sobolev space, Harnack inequality, $\phi$-Laplacian.

Mathematics Subject Classification: 35B50, 35J20, 35J60.

## 1. INTRODUCTION

Consider the following $\phi$-Laplacian equation

$$
\begin{equation*}
-\operatorname{div}\left(\phi(|\nabla u|) \frac{\nabla u}{|\nabla u|}\right)=\mathcal{B}(\cdot, u) \text { in } \Omega \tag{1.1}
\end{equation*}
$$

on the Orlicz-Sobolev space $W^{1} L_{\Phi}(\Omega)$, where $\Omega \subseteq \mathbb{R}^{N}$ is a bounded domain which has the segment property [5]. (Complete treatments and characterizations of the Orlicz-Sobolev spaces $W^{1} L_{\Phi}(\Omega)$ and $W_{0}^{1} L_{\Phi}(\Omega)$ can be found in $\left.[1,8,10]\right)$. The term $\phi: \mathbb{R} \rightarrow \mathbb{R}$ is an odd and increasing homeomorphism. However, in striking contrast with the classic case treated by Lieberman [11], $\phi$ is not required to be differentiable. It is rather assumed that

$$
\begin{equation*}
1<p_{\Phi} \leq \frac{t \phi(t)}{\Phi(t)} \leq q_{\Phi}<+\infty, \quad t \neq 0 \tag{1.2}
\end{equation*}
$$

for some numbers $p_{\Phi}$ and $q_{\Phi}$, where $\Phi$ is the $N$-function

$$
\begin{equation*}
\Phi(t)=\int_{0}^{t} \phi(s) d s \tag{1.3}
\end{equation*}
$$

Note that integration of eq. (1.2) yields $\Phi(\kappa t) \leq \kappa^{q_{\Phi}} \Phi(t)$ for $\kappa>1$ and $t>0$ and this implies the useful inequality

$$
\begin{equation*}
\phi(\kappa t) \leq q_{\Phi} \kappa^{q_{\Phi}-1} \phi(t) \quad \text { for all } t \geq 0 \text { and } \kappa>1 \tag{1.4}
\end{equation*}
$$

The right-hand side $\mathcal{B}: \Omega \times W^{1} L_{\Phi}(\Omega) \rightarrow \mathbb{R}$ of (1.1) is a Carathéodory function such that

$$
\begin{equation*}
|\mathcal{B}(x, u)| \leq \mathfrak{a} \phi(|u(x)|)+\mathfrak{b} \text { a.e. } x \text { in } \Omega, \tag{1.5}
\end{equation*}
$$

where $\mathfrak{a}, \mathfrak{b}$ are two nonnegative numbers. Let $B_{R} \subset \subset \Omega$ be a ball of radius $0<R \leq 1$ and let $B_{R / 2}$ be the concentric ball of radius $R / 2$. It was proved in [2] that if $u$ is a locally bounded and nonnegative solution of (1.1) then

$$
\sup _{B_{R / 2}} u \leq \mathcal{N}\left(\inf _{B_{R / 2}} u+L R\right)
$$

where $\mathcal{N}=\mathcal{N}\left(\mathfrak{a}, p_{\Phi}, q_{\Phi}, N\right)$ is a positive constant and $L>0$ is any constant such that $\mathfrak{b} \leq \phi(L)$. In this short note we present a parametric version of the estimates obtained in [2]. In this sense, the bounds obtained here are optimal. These improved bounds allow for a more general interpretation of the behavior of solutions of $\phi$-Laplacians and permit us to treat the problem on the regularity of the solutions [3]. Even though we do not address these properties here, it is of particular interest the special case of the $p$-Laplacian operator for which $\phi(s)=|s|^{p-2} s$ and $p>1$. The particular case of variable exponents $p(x)$, where $p: \Omega \rightarrow(1,+\infty)$ is a bounded function, is treated in [14-16]. Special types of nonlinearities in connection with the $p$-Laplace operator have been considered recently in the article [9]. In a rather different context, the article [12] provides a geometric approach to the study of the $p$-Laplacian on a ball in $\mathbb{R}^{N}$ using techniques from dynamical systems. The author studies the invariant manifolds at the union of the solutions in the phase space and addresses the variational aspects of the corresponding tangent vector fields.

## 2. IMPROVED ESTIMATES

In this article, a solution of (1.1) will be any function $u \in W^{1} L_{\Phi}(\Omega)$ which satisfies estimates (1.5) and fulfills the identity

$$
\begin{equation*}
\int_{\Omega} \phi(|\nabla u|) \frac{\nabla u}{|\nabla u|} \cdot \nabla \theta d x=\int_{\Omega} \mathcal{B}(x, u) \theta d x \text { for all } \theta \in W_{0}^{1} L_{\Phi}(\Omega) . \tag{2.1}
\end{equation*}
$$

(Hölder's inequality guarantees that the integrals above are finite [3]).
Let $B_{s} \subseteq B_{R}$ be two concentric balls of radii $0<s<R \leq 1$ centered at a given point $x_{0} \in \mathbb{R}^{N}$. A smooth function $\eta: \mathbb{R}^{N} \rightarrow[0,1]$ satisfying simultaneously

$$
\left.\eta\right|_{B_{s}} \equiv 1,\left.\quad \eta\right|_{\left(\mathbb{R}^{N} \backslash B_{R}\right)} \equiv 0, \quad|\nabla \eta| \leq \frac{2}{R-s}
$$

is called an $s$-cut-off function on $B_{R}$. Such functions do exist [7].

Lemma 2.1 ([2, Lemma 3.3]). Let $B_{R} \subset \subset \Omega$ be any ball of radius $0<R \leq 1$ and let $\eta$ be an $R / 2$-cut-off function on $B_{R}$. Choose $v \geq 0$ such that $v^{\alpha} \in L^{\infty}\left(B_{R}\right)$, where $\alpha= \pm 1$. Set $w=R^{-1} \eta^{\alpha} v$. Then $q_{\Phi}{ }^{-1} \sup _{B_{R}} w^{\alpha} \sup _{B_{R}}(\phi(w))^{\alpha} \leq \sup _{B_{R}}(\Phi(w))^{\alpha}$.
Proposition 2.2. Let $B_{R} \subset \subset \Omega$ be any ball of radius $0<R \leq 1$. Choose a $\sigma R$-cut-off function $\eta$ on $B_{R}$, where $\sigma \in(0,1)$. Suppose that $v$ is a nonnegative function such that $v^{\alpha} \in L^{\infty}\left(B_{R}\right)$, where $\alpha= \pm 1$. Set $w=R^{-1} \eta^{\alpha} v$ and assume that

$$
\begin{equation*}
\left(\int_{B_{R}} \eta^{-\beta}(\Phi(w))^{\alpha q Q} d x\right)^{1 / Q} \leq \frac{C q}{R(1-\sigma)^{q_{\Phi}}} \int_{B_{R}} \eta^{-\beta}(\Phi(w))^{\alpha q} d x \tag{2.2}
\end{equation*}
$$

where the numbers $q \geq \beta>0, Q=N /(N-1)$ and the constant $C$ depends neither on the ball $B_{R}$ nor on the number $q$. Then for $p>0$ there exists a positive constant $\mathcal{C}=\mathcal{C}\left(\beta, p, q_{\Phi}, C, N\right)$ such that

$$
\sup _{B_{\sigma R}} v^{\alpha} \leq \frac{\mathcal{C}}{(1-\sigma)^{N q_{\Phi} / p}}\left(f_{B_{R}} v^{\alpha p} d x\right)^{1 / p}
$$

where $B_{\sigma R}$ is the ball of radius $\sigma R$ concentric with $B_{R}$.
Proof. Assume first that $\beta \leq p$. The particular choice $q=p Q^{\nu}$ in (2.2), where $\nu$ is a nonnegative integer, produces

$$
\begin{aligned}
& \left(\int_{B_{R}} \eta^{-\beta}(\Phi(w))^{\alpha p Q^{\nu+1}} d x\right)^{1 / p Q^{\nu+1}} \\
& \leq \frac{\left(C p Q^{\nu}\right)^{1 / p Q^{\nu}}}{R^{1 / p Q^{\nu}}(1-\sigma)^{q_{\Phi} / p Q^{\nu}}}\left(\int_{B_{R}} \eta^{-\beta}(\Phi(w))^{\alpha p Q^{\nu}} d x\right)^{1 / p Q^{\nu}}
\end{aligned}
$$

For $m \geq 1$, consider $\nu$ large enough such that $p Q^{\nu+1}>m$. Since $(\Phi(w))^{\alpha} \in L^{\infty}\left(B_{R}\right)$, a Moser iteration [13] of this inequality with respect to $\nu$ and the imbedding theorem yield

$$
\left\|(\Phi(w))^{\alpha}\right\|_{L^{m}\left(B_{R}\right)} \leq\left(\frac{C p}{R(1-\sigma)^{q_{\Phi}}}\right)^{N / p} Q^{N(N-1) / p}\left(\int_{B_{R}} \eta^{-\beta}(\Phi(w))^{\alpha p} d x\right)^{1 / p}
$$

where $Q=N /(N-1)$. This estimate is valid for the norm $L^{\infty}(\Omega)$ [1, Theorem 2.14]. Since $(\Phi(w))^{\alpha} \leq q_{\Phi} w^{\alpha}(\phi(w))^{\alpha}$ and $w^{\alpha}=R^{-\alpha} \eta^{\alpha^{2}} v^{\alpha}$, Lemma 2.1 implies

$$
\begin{equation*}
\sup _{B_{\sigma R}} v^{\alpha} \leq q_{\Phi}^{2}\left(\frac{C p}{(1-\sigma)^{q_{\Phi}}}\right)^{N / p} Q^{N(N-1) / p} \omega_{N}^{1 / p}\left(f_{B_{R}} v^{\alpha p} d x\right)^{1 / p} \tag{2.3}
\end{equation*}
$$

where $\omega_{N}$ denotes the volume of the unit ball in $\mathbb{R}^{N}$. This concludes the proof in the case $p \geq \beta$, with $\mathcal{C}=\mathcal{D}_{1}(p):=q_{\Phi}^{2}(C p)^{N / p} Q^{N(N-1) / p} \omega_{N}^{1 / p}$. Note that if we write $s=\sigma R$ then the previous estimate, applied with $p=\beta$, reads

$$
\begin{equation*}
\sup _{B_{s}} v^{\alpha} \leq \frac{\mathcal{D}_{1}(\beta)}{(1-s / R)^{N q_{\Phi} / \beta}}\left(f_{B_{R}} v^{\alpha \beta} d x\right)^{1 / \beta} \tag{2.4}
\end{equation*}
$$

Now, suppose that $0<p<\beta$. Since $v^{\alpha \beta}=v^{\alpha(\beta-p)} v^{\alpha p}$, estimate (2.4) produces

$$
\sup _{B_{s}} v^{\alpha} \leq \frac{\mathcal{D}_{1}(\beta)}{(1-s / R)^{N q_{\Phi} / \beta}}\left(\sup _{B_{R}} v^{\alpha}\right)^{1-p / \beta}\left(f_{B_{R}} v^{\alpha p} d x\right)^{1 / \beta}
$$

Since $c d \leq \frac{(\beta-p)}{\beta} c^{\frac{\beta}{(\beta-p)}}+\frac{p}{\beta} d^{\frac{\beta}{p}}$ for $c, d \geq 0$, the inequality below follows:

$$
\sup _{B_{s}} v^{\alpha} \leq \frac{(\beta-p)}{\beta} \sup _{B_{R}} v^{\alpha}+\frac{p}{\beta} \frac{\mathcal{D}_{1}{ }^{\beta / p}}{(R-s)^{N q_{\Phi} / p}}\left(R^{N q_{\Phi}} \int_{B_{R}} v^{\alpha p} d x\right)^{1 / p}
$$

An application of $[6$, Lemma 3.1] on the interval $[0, R]$ yields the conclusion.
Let $L$ be any real nonnegative constant such that $\mathfrak{b} \leq \phi(L)$. Choose $0<R \leq 1$ and write $v=u+R L$ where $u$ is a nonnegative solution of (1.1). Then

$$
\mathfrak{a} \phi(u)+\mathfrak{b} \leq \mathfrak{a} \phi(u)+\phi(L) \leq \mathfrak{a} \phi(u+R L)+\phi\left(\frac{R L}{R}\right) \leq \mathfrak{a} \phi\left(\frac{R v}{R}\right)+\phi\left(\frac{v}{R}\right)
$$

By (1.5), the right-hand side of the equation (1.1) is hence bounded as follows:

$$
\begin{equation*}
|\mathcal{B}(\cdot, u)| \leq M \phi\left(\frac{v}{R}\right) \tag{2.5}
\end{equation*}
$$

with the constant $M=\mathfrak{a}+1$.
Proposition 2.3 ([2, Proposition 2.1]). Let $u \in W^{1} L_{\Phi}(\Omega)$. If $\operatorname{supp}(u)$ is compact contained in $\Omega$ then $u \in W_{0}^{1} L_{\Phi}(\Omega)$.
Proposition 2.4. Let $B_{R} \subset \subset \Omega$ be any ball of radius $0<R \leq 1$. Choose a $\sigma R$-cut-off function $\eta$ on $B_{R}$ where $\sigma \in(0,1)$. Suppose that $u$ is a locally bounded and nonnegative solution of eq. (1.1). Choose any $L \geq 0$ such that $\mathfrak{b} \leq \phi(L)$. Then for any $d>0$ and $q \geq 2+d$,

$$
\begin{equation*}
\int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u| d x \leq \frac{C_{1}}{(1-\sigma)^{q_{\Phi}}} \int_{\Omega} \eta^{-d-2 q_{\Phi}}(\Phi(w))^{q} d x \tag{2.6}
\end{equation*}
$$

where $w=\eta(u+R L) / R$ and $C_{1}=C_{1}\left(\mathfrak{a}, q_{\Phi}\right)$ is a positive constant.

Proof. Define $v=u+R L$. The following standard argument will be repeatedly employed. Since $\operatorname{supp}(\eta)$ is compact in $\Omega$ the function $w=\eta v / R \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and thus there exists a constant $\mathrm{A}>0$ such that $|w(x)| \leq \mathrm{A}$ for all $x \in \mathbb{R}^{N}$. Define for all $t \geq 0$ the function $f_{1}(t)=t^{-d}(\Phi(t))^{q-1}$ if $t \in(0, \mathrm{~A}+1)$ and write $f_{1}(t)=0$ if $t \notin(0, \mathrm{~A}+1)$. Note that $t \mapsto t^{-d}(\Phi(t))^{q-1}$ is of class $C^{1}(0,+\infty)$ and the derivative of this map tends to zero as $t \rightarrow 0^{+}$. Hence, $f_{1}^{\prime}$ is uniformly bounded on $[0,+\infty)$. If $F_{1}$ denotes the odd extension of $f_{1}$ to the entire real line then $F_{1} \in C^{1}(\mathbb{R})$ and $F_{1}^{\prime} \in L^{\infty}(\mathbb{R})$. Next, define the term $\theta=R^{-d} v^{d+1} F_{1}(w)=\eta^{-d}(\Phi(w))^{q-1} v$. The product formula [7, eq. (7.18)] and [7, Theorem 7.8] yield

$$
\nabla \theta=\frac{(d+1)}{R^{d}} v^{d} F_{1}(w) \nabla u+\frac{v^{d+1}}{R^{d+1}} F_{1}^{\prime}(w) v \nabla \eta+\frac{v^{d+1}}{R^{d+1}} F_{1}^{\prime}(w) \eta \nabla u \in L_{\Phi}(\Omega)
$$

Since $\operatorname{supp}(\theta) \subseteq \bar{B}_{R} \subseteq \Omega$, Proposition 2.3 implies $\theta \in W_{0}^{1} L_{\Phi}(\Omega)$ and eq. (2.1) becomes

$$
\begin{aligned}
& \int_{\Omega} \eta^{-d}\left[(q-1)(\Phi(w))^{q-2} \phi(w) w+(\Phi(w))^{q-1}\right] \phi(|\nabla u|)|\nabla u| \\
& =\int_{\Omega} \mathcal{B}(x, u) \eta^{-d}(\Phi(w))^{q-1} v \\
& \quad-(q-1) R \int_{\Omega} \eta^{-d-2}(\Phi(w))^{q-2} \phi(w) w^{2} \phi(|\nabla u|) \frac{\nabla u}{|\nabla u|} \nabla \eta \\
& \quad+d R \int_{\Omega} \eta^{-d-2}(\Phi(w))^{q-1} w \phi(|\nabla u|) \frac{\nabla u}{|\nabla u|} \nabla \eta .
\end{aligned}
$$

The term $\phi(w) w$ in the argument of the integral on the left-hand side of the equality is bounded from below by $\Phi(w)$. The absolute value of the integrals on the right is taken. Then the bound $\phi(w) w \leq q_{\Phi} \Phi(w)$ is applied to the argument of the second integral on the right. Since $q \geq \max \{1, d\}$, bounds $|\nabla \eta| \leq 2 / R(1-\sigma)$ and (2.5) produce

$$
\begin{align*}
& q \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u| d x \leq M q \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi\left(\frac{v}{R}\right) v d x \\
& \quad+4 q q_{\Phi} \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \frac{w}{\eta^{2}} \frac{\phi(|\nabla u|)}{(1-\sigma)} d x+2 q \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \frac{w}{\eta^{2}} \frac{\phi(|\nabla u|)}{(1-\sigma)} d x  \tag{2.7}\\
& =M q \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi\left(\frac{v}{R}\right) v d x+\mu q \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \frac{w}{\eta^{2}} \frac{\phi(|\nabla u|)}{(1-\sigma)} d x
\end{align*}
$$

where $\mu=2\left(2 q_{\Phi}+1\right)$. Division by $q$ is performed on both sides of the inequality. The second integral on the right is multiplied and divided by a sufficiently small quantity $\varepsilon>0$ so as to form the term $w / \varepsilon \eta^{2}$ in the argument. Young's inequality
$t \phi(s) \leq t \phi(t)+s \phi(s)$, with $t=w / \varepsilon \eta^{2}(1-\sigma)$ and $s=|\nabla u|$, is applied to this integral.
From eq. (2.7) we obtain

$$
\begin{aligned}
& (1-\varepsilon \mu) \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u| d x \\
& \leq M \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi\left(\frac{v}{R}\right) v d x+\mu \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \frac{w}{\eta^{2}(1-\sigma)} \phi\left(\frac{w}{\varepsilon \eta^{2}(1-\sigma)}\right) d x
\end{aligned}
$$

Eq. (1.4) yields the following two estimates:

$$
\phi(v / R) v=\phi(w / \eta) R w / \eta \leq q_{\Phi} \eta^{-q_{\Phi}} \phi(w) w
$$

and

$$
\phi\left(w / \varepsilon \eta^{2}(1-\sigma)\right) \leq q_{\Phi} \phi(w) /\left(\varepsilon \eta^{2}(1-\sigma)\right)^{q_{\Phi}-1}
$$

Hence

$$
\begin{aligned}
&(1-\varepsilon \mu) \int_{\Omega} \eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u| d x \\
& \leq M q_{\Phi} \int_{\Omega} \eta^{-d-q_{\Phi}}(\Phi(w))^{q-1} \phi(w) w d x \\
& \quad+\frac{\mu q_{\Phi}}{\varepsilon^{q_{\Phi}-1}(1-\sigma)^{q_{\Phi}}} \int_{\Omega} \eta^{-d-2 q_{\Phi}}(\Phi(w))^{q-1} \phi(w) w d x
\end{aligned}
$$

For $\varepsilon$ suitably chosen (e.g. $\varepsilon=1 / 2 \mu$ ) and since $\phi(w) w \leq q_{\Phi} \Phi(w)$ and $\eta^{-q_{\Phi}} \leq \eta^{-2 q_{\Phi}}$, the conclusion follows.

Corollary 2.5. Let $u$ be a locally bounded and nonnegative solution of equation (1.1) and let $B_{R} \subset \subset \Omega$ be a ball of radius $0<R \leq 1$. Then for any $p>0$,

$$
\sup _{B_{\sigma R}} u \leq \frac{2^{1+1 / p} \mathcal{C}}{(1-\sigma)^{N q_{\Phi} / p}}\left(\left(f_{B_{R}} u^{p} d x\right)^{1 / p}+R L\right)
$$

where $\mathcal{C}=\mathcal{C}\left(\mathfrak{a}, p, q_{\Phi}, N\right)$ is the constant in Proposition 2.2 and $\sigma \in(0,1)$. The term $L$ is any real nonnegative constant such that $\mathfrak{b} \leq \phi(L)$ and $B_{\sigma R}$ is the ball of radius $\sigma R$ concentric with $B_{R}$.

Proof. Let $\eta$ be an $\sigma R$-cut-off function on $B_{R}$. As in the proof of Proposition 2.4, if we set $v=u+R L$ then $w=\eta v / R=\eta(u+R L) / R \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and thus $|w(x)| \leq \mathrm{A}$ for $x \in \mathbb{R}^{N}$ and where A is the same constant in the proof of Proposition 2.4. Take $d>0, q \geq 2+d$ (to be fixed later) and define $f_{2}(t)=\Phi(t) f_{1}(t)$ for $t \geq 0$ and where $f_{1}$ is the function defined in the proof of Proposition 2.4. It is evident that $f_{2}^{\prime}$ is bounded
on $[0,+\infty)$. It follows that $F_{2} \in C^{1}(\mathbb{R})$ and $F_{2}^{\prime} \in L^{\infty}(\mathbb{R})$ where $F_{2}$ is the odd extension of $f_{2}$ to the entire real line. Define $\theta=F_{2}(w) v^{d} / R^{d}=\eta^{-d}(\Phi(w))^{q}$. Theorem 7.8 in [7] yields

$$
\begin{aligned}
\nabla \theta= & -d \eta^{-d-1}(\Phi(w))^{q} \nabla \eta+\eta^{-d-1} q(\Phi(w))^{q-1} \phi(w) w \nabla \eta \\
& +\frac{q}{R} \eta^{-d+1}(\Phi(w))^{q-1} \phi(w) \nabla u
\end{aligned}
$$

Note that $\phi(w) w \leq q_{\Phi} \Phi(w), q \geq d$ and $|\nabla \eta| \leq 2 / R(1-\sigma)$. Then Gagliardo-Nirenberg inequality [4, Theorem IX], applied to $\theta$ and $Q=N /(N-1)$, yields $\mathcal{C}_{N}>0$ such that

$$
\begin{aligned}
& \left(\int_{\Omega}\left|\eta^{-d}(\Phi(w))^{q}\right|^{Q}\right)^{1 / Q} \\
& \leq \mathcal{C}_{N}\left(\frac{2 q\left(1+q_{\Phi}\right)}{R(1-\sigma)} \int_{\Omega} \eta^{-d-1}(\Phi(w))^{q}+\frac{q}{R} \int_{\Omega} \eta^{-d+1}(\Phi(w))^{q-1} \phi(w)|\nabla u|\right) .
\end{aligned}
$$

Since $\eta \leq 1$ the first integral on the right-hand side of the inequality above is bounded by $\int_{\Omega} \eta^{-d-2 q_{\Phi}}(\Phi(w))^{q} d x$. The argument of the second integral on the right is bounded as follows. Estimate $t \phi(s) \leq t \phi(t)+s \phi(s)$ is again used (with $t=|\nabla u|$ and $s=w$ ) and hence

$$
\begin{aligned}
\eta^{-d+1}(\Phi(w))^{q-1} \phi(w)|\nabla u| & \leq \eta^{-d}(\Phi(w))^{q-1} \phi(w) w+\eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u| \\
& \leq q_{\Phi} \eta^{-d-2 q_{\Phi}}(\Phi(w))^{q}+\eta^{-d}(\Phi(w))^{q-1} \phi(|\nabla u|)|\nabla u|
\end{aligned}
$$

Since $1-\sigma<1$, eq. (2.6) and the choice $-d Q=-d-2 q_{\Phi}$ (i.e. $d=2(N-1) q_{\Phi}$ ) yield

$$
\left(\int_{B_{R}} \eta^{-2 N q_{\Phi}}(\Phi(w))^{q Q} d x\right)^{1 / Q} \leq \frac{C_{2} q}{R(1-\sigma)^{q_{\Phi}}} \int_{B_{R}} \eta^{-2 N q_{\Phi}}(\Phi(w))^{q} d x
$$

where $C_{2}=C_{2}\left(\mathfrak{a}, q_{\Phi}, N\right)=\mathcal{C}_{N}\left(5 q_{\Phi}+C_{1}\right)$. Proposition 2.2 with $C=C_{2}, \beta=2 N q_{\Phi}$, $p>0$ and $\alpha=1$ gives

$$
\sup _{B_{\sigma R}} u \leq \sup _{B_{\sigma R}} v \leq \frac{\mathcal{C}}{(1-\sigma)^{N q_{\Phi} / p}}\left(f_{B_{R}} v^{p} d x\right)^{1 / p}
$$

for the same constant $\mathcal{C}$ in that proposition (and which, in this case, clearly depends on $\mathfrak{a}, p, q_{\phi}$ and $\left.N\right)$. Note that $x \mapsto x^{r}$ is convex in $(0,+\infty)$ for $r \geq 1$ whereas $x \mapsto 1+x^{r}-(1+x)^{r}$ is increasing in $(0,+\infty)$ if $0<r<1$. Hence $(c+d)^{r} \leq 2^{r}\left(c^{r}+d^{r}\right)$ for $c, d, r>0$ and this implies the bound claimed. The corollary is proved.

Proposition 2.6. Let $B_{R} \subset \subset \Omega$ be any ball of radius $0<R \leq 1$. Choose a $\sigma R$-cut-off function $\eta$ on $B_{R}$ where $\sigma \in(0,1)$. Suppose that $u$ is a locally bounded and nonnegative solution of eq. (1.1). Choose any $L \geq 0$ such that $\mathfrak{b} \leq \phi(L)$ and let $\chi>0$ be arbitrary. Then for any $q \geq d \geq 1$,

$$
\begin{equation*}
\int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} \phi(|\nabla u|)|\nabla u| d x \leq \frac{C_{3}}{(1-\sigma)^{q_{\Phi}}} \int_{\Omega} \eta^{-d}(\Phi(w))^{-q} d x \tag{2.8}
\end{equation*}
$$

where $w=(u+R L+\chi) / R \eta$ and $C_{3}=C_{3}\left(\mathfrak{a}, q_{\Phi}\right)$ is a positive constant.
Proof. Define $v=u+R L$ and $z=R \eta / \psi$, where $\psi=v+\chi=u+R L+\chi$. Since $u$ is nonnegative, $z \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and there exists $\mathrm{B}>0$ such that $|z(x)| \leq \mathrm{B}$ for all $x \in \mathbb{R}^{N}$. Consider for all $t \geq 0$ the function $f_{3}(t)=t^{-d}\left(\Phi\left(t^{-1}\right)\right)^{-(q+1)}$ if $t \in(0, \mathrm{~B}+1)$ and $f_{3}(t)=0$ if $t \notin(0, \mathrm{~B}+1)$. In this case, the map $t \mapsto t^{-d}\left(\Phi\left(t^{-1}\right)\right)^{-(q+1)}$ is of class $C^{1}(0, \infty)$ and its derivative tends to zero as $t \rightarrow 0^{+}$. It follows that $f_{3}^{\prime}$ is uniformly bounded on $[0,+\infty)$. If $F_{3}$ is the odd extension of $f_{3}$ to the entire real line then $F_{3} \in C^{1}(\mathbb{R})$ and $F_{3}^{\prime} \in L^{\infty}(\mathbb{R})$. Note that $w=z^{-1}$. Define

$$
\theta=R^{d} F_{3}(z) \psi^{1-d}=R \eta^{1-d}(\Phi(w))^{-q-1} w
$$

In this case,

$$
\begin{aligned}
\nabla \theta= & R^{d+1} F_{3}^{\prime}(z) \psi^{-d} \nabla \eta-R^{d+1} \eta F_{3}^{\prime}(z) \psi^{-(d+1)} \nabla u \\
& +(1-d) R^{d} F_{3}(z) \psi^{-d} \nabla u \in L_{\Phi}(\Omega)
\end{aligned}
$$

Since $\operatorname{supp}(\theta) \subseteq \bar{B}_{R} \subseteq \Omega$, Proposition 2.3 ensures that the function $\theta \in W_{0}^{1} L_{\Phi}(\Omega)$. The weak formulation (2.1) in this case is equivalent to

$$
\begin{aligned}
& \int_{\Omega} \eta^{-d}\left[(q+1)(\Phi(w))^{-q-2} \phi(w) w-(\Phi(w))^{-q-1}\right] \phi(|\nabla u|)|\nabla u| d x \\
& =R(q+1) \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-2} \phi(w) w^{2} \phi(|\nabla u|) \frac{\nabla u}{|\nabla u|} \nabla \eta d x \\
& \quad-R d \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} w \phi(|\nabla u|) \frac{\nabla u}{|\nabla u|} \nabla \eta d x \\
& \quad-R \int_{\Omega} \mathcal{B}(x, u) \eta^{1-d}(\Phi(w))^{-q-1} w d x
\end{aligned}
$$

The estimate $\Phi(w) \leq \phi(w) w$ is applied to the argument of the integral on the left-hand side of the equality. After taking absolute values on the right-hand side, the bound $\phi(w) w \leq q_{\Phi} \Phi(w)$ is applied to the argument of the first integral on the right-hand side.

Write $\mu=2\left(2 q_{\Phi}+1\right)$. Since $q \geq d \geq 1$ and $|\nabla \eta| \leq 2 / R(1-\sigma)$ the above equivalence yields

$$
\begin{aligned}
& q \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} \phi(|\nabla u|)|\nabla u| d x \\
& \leq q \mu \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} \frac{w}{(1-\sigma)} \phi(|\nabla u|) d x+M q \int_{\Omega} \phi\left(\frac{v}{R}\right) \eta^{1-d}(\Phi(w))^{-q-1} w d x .
\end{aligned}
$$

The term $w$ in the argument of the first integral on the right is divided and multiplied $\varepsilon>0$ small. Equation (1.4) implies

$$
\phi(w / \varepsilon(1-\sigma)) \leq q_{\Phi}(\varepsilon(1-\sigma))^{-q_{\Phi}+1} \phi(w)
$$

Since $\phi(v / R) \leq \phi(\eta w) \leq \phi(w)$ and $\eta^{1-d} \leq \eta^{-d}$, the bound $t \phi(s) \leq t \phi(t)+s \phi(s)$ with $t=w / \varepsilon(1-\sigma)$ and $s=|\nabla u|$ yields

$$
\begin{aligned}
& (1-\varepsilon \mu) \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} \phi(|\nabla u|)|\nabla u| \\
& \leq \frac{q_{\Phi}^{2} \mu}{\varepsilon^{q_{\Phi}-1}(1-\sigma)^{q_{\Phi}}} \int_{\Omega} \eta^{-d}(\Phi(w))^{-q}+M q_{\Phi} \int_{\Omega} \eta^{-d}(\Phi(w))^{-q} .
\end{aligned}
$$

If $\varepsilon$ is suitably chosen then the conclusion follows.
Corollary 2.7 (Weak Harnack inequality). Let u be a locally bounded and nonnegative solution of equation (1.1). Let $B_{R} \subset \subset \Omega$ be a ball of radius $0<R \leq 1$ and $\sigma \in(0,1)$. Then there exist positive constants $p_{0}=p_{0}\left(\mathfrak{a}, p_{\Phi}, q_{\Phi}, N\right)$ and $\mathrm{C}=\mathrm{C}\left(\mathfrak{a}, p_{\Phi}, q_{\Phi}, N\right)$ such that

$$
\left(f_{B_{R}} u^{p_{0}} d x\right)^{1 / p_{0}} \leq \frac{\mathrm{C}}{(1-\sigma)^{N q_{\Phi} / p_{0}}}\left(\inf _{B_{\sigma R}} u+R L\right)
$$

where $L$ is any real nonnegative constant such that $\mathfrak{b} \leq \phi(L)$ and $B_{\sigma R}$ is the ball of radius $\sigma R$ concentric with $B_{R}$.

Proof. Let $\eta$ be an $\sigma R$-cut-off function on $B_{R}$ and take $\chi>0$, arbitrary. As in the proof of Proposition 2.6, $z:=R \eta /(u+R L+\chi) \in L^{\infty}\left(\mathbb{R}^{N}\right)$ and thus $|z(x)| \leq \mathrm{B}$ for $x \in \mathbb{R}^{N}$ and for the same constant. Take $q \geq d \geq 1$ (to be fixed later) and for all $t \geq 0$ define $f_{4}(t)=t R^{d-1} \Phi\left(t^{-1}\right) f_{3}(t)$ which is differentiable with bounded derivative. It is clear that if $F_{4}$ is the odd extension of $f_{4}$ to the entire real line then $F_{4} \in C^{1}(\mathbb{R})$ and $F_{4}^{\prime} \in L^{\infty}(\mathbb{R})$. Write $w=z^{-1}$. Then [7, Theorem 7.8] applied to $F_{4}(z) \psi^{1-d}=\eta^{1-d}(\Phi(w))^{-q}$ yields

$$
\begin{aligned}
& \nabla\left(\eta^{1-d}(\Phi(w))^{-q}\right) \\
& =\left[(1-d) \eta^{-d}(\Phi(w))^{-q}+q \eta^{-d}(\Phi(w))^{-q-1} \phi(w) w\right] \nabla \eta-\frac{q}{R} \eta^{-d}(\Phi(w))^{-q-1} \phi(w) \nabla u .
\end{aligned}
$$

Since $q \geq d \geq 1$, the bound $\phi(t) t \leq q_{\Phi} \Phi(t)$ and the Gagliardo-Nirenberg inequality yield

$$
\begin{aligned}
& \left(\int_{\Omega}\left|\eta^{1-d}(\Phi(w))^{-q}\right|^{Q}\right)^{1 / Q} \\
& \leq \frac{2 q \mathcal{C}_{N}\left(q_{\Phi}+2\right)}{R(1-\sigma)} \int_{\Omega} \eta^{-d}(\Phi(w))^{-q}+\frac{\mathcal{C}_{N} q}{R} \int_{\Omega} \eta^{-d}(\Phi(w))^{-q-1} \phi(w)|\nabla u|
\end{aligned}
$$

where $Q=N /(N-1)$. It follows from eq. (2.8) (with $d=N$ ) and Young's inequality that

$$
\left(\int_{B_{R}} \eta^{-N}(\Phi(w))^{-q Q} d x\right)^{1 / Q} \leq \frac{C_{4} q}{R(1-\sigma)^{q_{\Phi}}} \int_{B_{R}} \eta^{-N}(\Phi(w))^{-q} d x
$$

where $C_{4}=C_{4}\left(\mathfrak{a}, q_{\Phi}, N\right)$. Proposition 2.2 with $\beta=N, v=\psi$ and $\alpha=-1$ yields

$$
\begin{equation*}
\left(f_{B_{R}} \psi^{-p} d x\right)^{-1} \leq \frac{\mathcal{C}^{p}}{(1-\sigma)^{N q_{\Phi}}}\left(\inf _{B_{R / 2}} \psi^{-1}\right)^{p} \tag{2.9}
\end{equation*}
$$

where $p>0$ and $\mathcal{C}=\mathcal{C}\left(\mathfrak{a}, p, q_{\Phi}, N\right)$. Let $\xi$ be a $\rho$-cut-off function on $B_{R}$ where $0<\rho \leq R / 2$. Consider the function $\theta=\xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \psi$, where $\mathbf{w}=\psi / \rho$. It is easy to prove that
$\nabla \theta=-\frac{1}{\rho}(\Phi(\mathbf{w}))^{-2} \phi(\mathbf{w}) \xi^{q_{\Phi}} \psi \nabla v+q_{\Phi}(\Phi(\mathbf{w}))^{-1} \xi^{q_{\Phi}-1} \psi \nabla \xi+(\Phi(\mathbf{w}))^{-1} \xi^{q_{\Phi}} \nabla v \in L_{\Phi}(\Omega)$ with $\nabla v=\nabla u$. Since $\theta \in W_{0}^{1} L_{\Phi}(\Omega)$, the bound $|\nabla \xi| \leq 4 /(R-\rho) \leq 4 / \rho$ along with eq. (2.1) and the obvious bound

$$
(\Phi(t))^{-2} \phi(t) t-(\Phi(t))^{-1} \geq(\Phi(t))^{-1}\left(p_{\Phi}-1\right)>0
$$

yield

$$
\begin{aligned}
& \left(p_{\Phi}-1\right) \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \phi(|\nabla u|)|\nabla u| \\
& \leq 4 \varepsilon q_{\Phi} \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \frac{\mathbf{w}}{\varepsilon \xi} \phi(|\nabla u|)+M \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \phi\left(\frac{v}{R}\right) \frac{\rho \psi}{\rho}
\end{aligned}
$$

where $\varepsilon>0$ is sufficiently small. By Young's inequality (with $t=\mathbf{w} / \varepsilon \xi$ and $s=|\nabla u|$ ),

$$
\begin{aligned}
& \left(p_{\Phi}-1-4 \varepsilon q_{\Phi}\right) \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \phi(|\nabla u|)|\nabla u| \\
& \leq 4 q_{\Phi} \int_{\Omega} \xi^{q_{\Phi}-1}(\Phi(\mathbf{w}))^{-1} \mathbf{w} \phi\left(\frac{\mathbf{w}}{\varepsilon \xi}\right)+M \rho \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \phi\left(\frac{v}{R}\right) \mathbf{w} .
\end{aligned}
$$

$\operatorname{By}(1.4), \phi(\mathbf{w} / \varepsilon \xi) \leq q_{\Phi} \phi(\mathbf{w}) /(\varepsilon \xi)^{q_{\Phi}-1}$ and since $\phi(v / R) \leq \phi(\psi / R) \leq \phi(\psi / \rho)=\phi(\mathbf{w})$,

$$
\left(p_{\Phi}-1-4 \varepsilon q_{\Phi}\right) \int_{\Omega} \xi^{q_{\Phi}}(\Phi(\mathbf{w}))^{-1} \phi(|\nabla u|)|\nabla u| d x \leq\left(\frac{4 q_{\Phi}^{3}}{\varepsilon^{q_{\Phi}-1}}+M q_{\Phi}\right) \int_{B_{\rho}} 1 d x
$$

If $\varepsilon$ is chosen sufficiently small such that $\varepsilon<\left(p_{\Phi}-1\right) / 4 q_{\Phi}$ then

$$
\int_{B_{\rho}} \xi^{q_{\Phi}} \frac{\phi(|\nabla u|)|\nabla u|}{\Phi(\mathbf{w})} d x \leq C_{5} \operatorname{vol}\left(B_{\rho}\right),
$$

where $\operatorname{vol}\left(B_{\rho}\right)$ is the Lebesgue measure of $B_{\rho}$ and $C_{5}=C_{5}\left(\mathfrak{a}, p_{\Phi}, q_{\Phi}\right)>0$. Since $|\nabla \psi| \phi(\mathbf{w}) \leq|\nabla u| \phi(|\nabla u|)+\phi(\mathbf{w}) \mathbf{w}$ and as $\Phi(\mathbf{w}) \leq \phi(\mathbf{w}) \mathbf{w}$ the following bound holds
$\int_{B_{\rho / 2}} \frac{|\nabla \psi|}{\psi} d x \leq \frac{1}{\rho}\left(\int_{B_{\rho / 2}} \frac{|\nabla u| \phi(|\nabla u|)}{\Phi(\mathbf{w})} d x+\int_{B_{\rho / 2}} 1 d x\right) \leq 2^{N-1} \omega_{N}\left(C_{5}+\frac{1}{2^{N}}\right)\left(\frac{\rho}{2}\right)^{N-1}$,
where $\omega_{N}$ is the Lebesgue measure of the unit ball in $\mathbb{R}^{N}$. The result [7, Theorem 7.21], with $z=\log (\psi)$ and $\Omega^{\prime}=B_{R}$, implies the existence of a positive constant $D_{N}$ such that

$$
\left(\int_{B_{R}} e^{p_{0} z} d x\right)\left(\int_{B_{R}} e^{-p_{0} z} d x\right) \leq D_{N}^{2}(2 R)^{2 N}
$$

where $p_{0}=p_{0}\left(\mathfrak{a}, p_{\Phi}, q_{\Phi}, N\right)=2 \sigma_{0} /\left(2^{2 N} C_{5}+2^{N}\right)$ and $\sigma_{0}=\sigma_{0}(N)$ is the constant produced by [7, Theorem 7.21]. Along with estimate (2.9) (with $p=p_{0}$ ) the latter yields

$$
\int_{B_{R}} \psi^{p_{0}} d x \leq \frac{\mathcal{C}^{p_{0}}}{(1-\sigma)^{N q_{\Phi}}}\left(\frac{D_{N} 2^{N}}{\omega_{N}}\right)^{2}\left(\inf _{B_{R / 2}} u+R L+\chi\right)^{p_{0}} .
$$

The conclusion follows after passing to the limit $\chi \rightarrow 0$. The proof is complete.
A combination of Corollary 2.5 and Corollary 2.7 produces the following improved version of the Harnack inequality.

Corollary 2.8 (Harnack inequality in arbitrary balls). Let u be a locally bounded and nonnegative solution of equation (1.1). Let $B_{R} \subset \subset \Omega$ be a ball of radius $0<R \leq 1$ and $\sigma \in(0,1)$. Then there exist positive constants $p_{0}$ and $\mathcal{M}$, which depend only on $\mathfrak{a}, p_{\Phi}, q_{\Phi}$ and $N$, such that

$$
\sup _{B_{\sigma R}} u \leq \frac{\mathcal{M}}{(1-\sigma)^{2 N q_{\Phi} / p_{0}}}\left(\inf _{B_{\sigma R}} u+L R\right),
$$

where $L$ is any real nonnegative constant such that $\mathfrak{b} \leq \phi(L)$ and $B_{\sigma R}$ is the ball of radius $\sigma R$ concentric with $B_{R}$.

## Acknowledgements

We are grateful to the referee for providing helpful observations and remarks.

## REFERENCES

[1] R.A. Adams, John J.J.F. Fournier, Sobolev Spaces, vol. 140, Pure and Applied Mathematics, 2nd ed., Elsevier B.V., Amsterdam, 2003.
[2] W. Arriagada, J. Huentutripay, A Harnack's inequality in Orlicz-Sobolev spaces, Studia Math. 243 (2018), 117-137.
[3] W. Arriagada, J. Huentutripay, Regularity, positivity and asymptotic vanishing of solutions of a $\phi$-Laplacian, Anal. Ştiinţ Univ. "Ovidius" Constanţa Ser. Mat. 25 (2017) 3, 59-72.
[4] H. Brezis, Analyse Fonctionnelle: Théorie et Applications, Masson, Paris, 1983.
[5] E. DiBenedetto, Partial Differential Equations, Cornerstones, Birkhäuser Boston, Inc., Boston, MA, 2nd ed., 2010.
[6] M. Giaquinta, Multiple Integrals in the Calculus of Variations and Nonlinear Elliptic System, Princeton Univ. Press, Princeton, New Jersey, 1983.
[7] D. Gilbarg, N. Trudinger, Elliptic Partial Differential Equations of Second Order, Springer-Verlag Germany, reprint of the 1998 edition, 2001.
[8] J.P. Gossez, Orlicz-Sobolev spaces and nonlinear elliptic boundary value problems, [in:] S. Fučík, A. Kufner (eds.), Nonlinear Analysis, Function Spaces and Applications, Proceedings of a Spring School held in Horní Bradlo, 1978, vol. 1, BSB B.G. Teubner Verlagsgesellschaft, Leipzig, 1979. Teubner Texte zur Mathematik, 59-94.
[9] U. Kaufmann, I. Medri, One-dimensional singular problems involving the p-Laplacian and nonlinearities indefinite in sign, Adv. Nonlinear Anal. 5 (2016), 251-259.
[10] M. Krasnosel'skii, J. Rutickii, Convex Functions and Orlicz Space, English translation P. Noordhoff Ltd., Groningen, 1961.
[11] G.M. Lieberman, The natural generalization of the natural conditions of Ladyzhenskaya and Ural'tseva for elliptic equations, Comm. Partial Differential Equations 16 (1991) 2-3, 311-361.
[12] B. Maultsby, Uniqueness of solutions to singular p-Laplacian equations with subcritical nonlinearity, Adv. Nonlinear Anal. 6 (2017) 1, 37-59.
[13] J. Moser, On Harnack's theorem for elliptic differential equations, Comm. Pure Appl. Math. 14 (1961), 577-591.
[14] V. Rădulescu, Nonlinear elliptic equations with variable exponent: old and new, Nonlinear Anal. 121 (2015), 336-369.
[15] V. Rădulescu, D. Repovš, Partial Differential Equations with Variable Exponents: Variational Methods and Qualitative Analysis, Chapman \& Hall/CRC Monographs and Research Notes in Mathematics, Taylor \& Francis Group, Boca Raton FL, 2015.
[16] I.-L. Stăncuţ, I.D. Stîrcu, Eigenvalue problems for anisotropic equations involving a potential on Orlicz-Sobolev type spaces, Opuscula Math. 36 (2016) 1, 81-101.

Waldo Arriagada
waldo.arriagada@ku.ac.ae
Khalifa University
Department of Applied Mathematics and Sciences
P.O. Box 127788, Abu Dhabi, United Arab Emirates

Jorge Huentutripay
jorge.huentutripay@uach.cl
Universidad Austral de Chile
Instituto de Ciencias Físicas y Matemáticas
Campus Isla Teja
Valdivia, Chile
Received: January 11, 2018.
Revised: February 21, 2018.
Accepted: February 21, 2018.

