

# Gold nano-elliptic arrays used in plasmonic waveguides in near infrared spectrum

TOFIQ NURMOHAMMADI<sup>1</sup>, REZA YADIPOUR<sup>1</sup>, KARIM ABBASIAN<sup>2\*</sup>

<sup>1</sup>University of Tabriz, Faculty of Electrical and Computer Engineering, Tabriz, Iran

<sup>2</sup>University of Tabriz, School of Engineering Emerging Technologies, Tabriz, Iran

\*Corresponding author: k\_abbasian@tabrizu.ac.ir

In this paper, we investigated plasmonic waveguides in near infrared spectrum, especially at original ( $\lambda = 1310$  nm) and also communication bands ( $\lambda = 1550$  nm) using gold nano-elliptic rings. It is possible to shift localized surface plasmon resonance, by appropriate geometrical properties, to the desired wavelength. Three-dimensional simulations utilizing the finite-difference time-domain algorithm are used to determine the set of geometrical parameters of gold nano-elliptic rings for exciting localized surface plasmon resonance at 1310 and 1550 nm. Employing different configurations of gold nano-elliptic rings chains, waveguides are designed, with  $-3$  dB transmission loss coefficients and group velocities calculations for different modes. In comparison with circular nanorings, elliptic rings showed better characteristics, such as high electric field enhancements and low loss transmission coefficients.

Keywords: loss coefficient, group velocity, finite-difference time-domain.

## 1. Introduction

Confinement of light in nanoscale structures, using noble metallic particles, is one of the targets in nanophotonics [1]. Miniaturization of optical integrated circuits, as one of the main aims, has been attended and progressed recently [2–4]. Plasmon waveguides are fundamental elements in optical integration technology because of highly localized field enhancement and enable guidance and confinement of electromagnetic waves under the diffraction limit [5]. The plasmonic properties of metallic nanostructures are dominated by the oscillations of free electrons of the nanostructure in resonance with the incident electromagnetic field. It is well known that surface plasmons polaritons (SPPs) could be generated inside the noble metal nanoparticles such as copper, gold, and silver [3, 6, 7]. The optical properties of remarked nanostructures are controlled by the coherent electron fluctuations in result of resonance during illumination by an incident light at a certain frequency, which is referred as localized surface

plasmon resonance (LSPR) [3, 6]. LSPR results in a strongly enhanced electric near field localized at the particle surface [7, 8], which has achieved tremendous attention because of its potential for various technologies such as field enhancement spectroscopy [9–12], nano-trapping [13], nonlinear optics [14] and optical waveguiding [15]. Recently, numerous researches in relation to coupled metallic nanostructures including silver cube clusters [16], gold necklaces [17] and gold particle chains [18], have been conducted in order to understand the plasmon coupling between nanometals.

A chain of nano-metal particles gives rise to very large field enhancements in their near region because of the near-field coupling of two adjacent particles. It is well known that LSPR oscillations are dependent on the geometry of the nano-metal, and then, by variation of geometrical parameters and nano-metal shape, the resonance wavelength and field enhancement could be engineered [19]. Recently, nanoring particles are especially noteworthy because of high local electric field generation by them as a result of powerful plasmon coupling of the outer and inner surfaces; therefore it is possible to tune plasmon frequency resonance from the visible to near-infrared (NIR) region by variation of the diameter and width of the ring [20, 21].

In this paper, plasmonic waveguides are investigated in the NIR region, especially at original (O-band,  $\lambda = 1310$  nm) and also communication bands (C-band,  $\lambda = 1550$  nm) using gold nano-elliptic rings, then they are compared with circular nanoring waveguides.

This paper is organized in the following way: in Section 2 we described gold nano-elliptic rings characteristics in the  $\text{SiO}_2$  host material; in Section 3, the nano-elliptic rings have been used in plasmonic waveguides and their characteristics have been obtained; in the last section, the paper is finalized by discussing the obtained results and conclusions.

## 2. Design of gold nano-elliptic rings

Host material of the proposed waveguide is  $\text{SiO}_2$ , where arrays of Au nano-elliptic rings are embedded. Plasmon resonance wavelength of the gold nanostructures is in

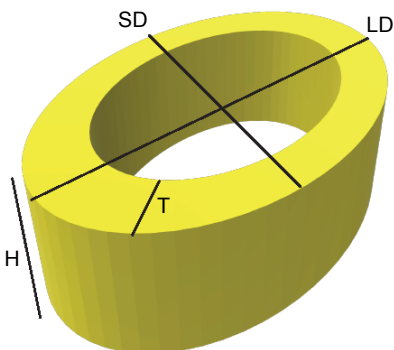


Fig. 1. Schematic illustration of nano-elliptic ring surrounded by  $\text{SiO}_2$ . The four independent geometrical parameters  $LD$ ,  $SD$ ,  $H$  and  $T$  are demonstrated ( $p = SD/LD$ ).

the range of  $\lambda = 520\text{--}600$  nm [22, 23], which is changed by varying geometrical parameters of nanostructures. Then the peak of their optical response is adjusted to the O-band or C-band. The object's geometrical degree of freedom provides the advantage of waveguide's performance optimizations. In nano-elliptic ring, the geometrical parameters are the elliptic large-diameter (LD), the elliptic small-diameter (SD), thickness  $T$  and height  $H$  (see Fig. 1); then we have four degrees of freedom in comparison with the cases of nanorings, where thickness, height and diameter are their degrees of freedom and sphere, where its radius is the only one degree of freedom.

Now, parameter  $p$  is introduced, where it is defined as the ratio of a SD to LD. In the sample, nano-elliptic ring parameters are  $LD = 235$  nm,  $H = 35$  nm, and  $T = 30$  nm. Figure 2 demonstrates LSPR blue shifting by decreasing  $p$ . On the other hand, when

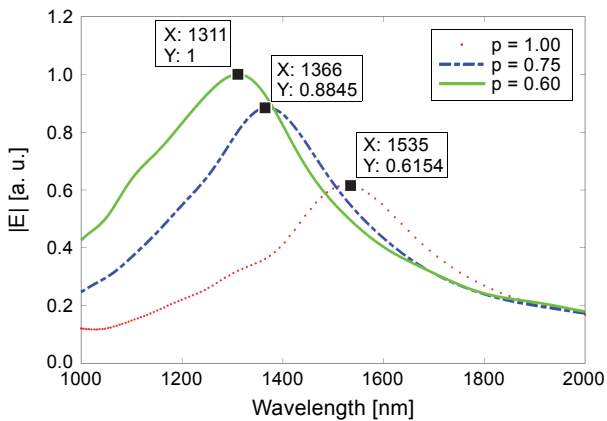


Fig. 2. Effect of parameter  $p$  on LSPR wavelength. By decreasing  $p$ , the LSPR is confronted to blue shift as well as higher electric field enhancement. The other geometrical parameters are  $LD = 235$  nm,  $H = 35$  nm, and  $T = 30$  nm.

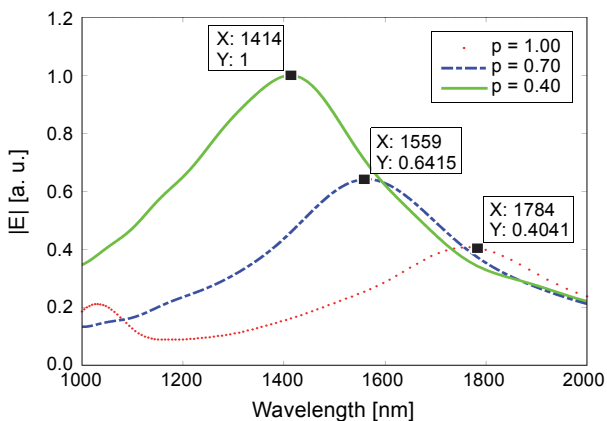


Fig. 3. Effect of parameter  $p$  on LSPR wavelength. By decreasing  $p$ , the LSPR is confronted to blue shift as well as higher electric field enhancement. The other geometrical parameters are  $LD = 290$  nm,  $H = 35$  nm, and  $T = 30$  nm.

T a b l e 1. FDTD simulation parameters.

Parameter	Values
Yee cell size ( $d_x = d_y = d_z$ )	4 nm
Simulation time	1000 fs
Boundary conditions	PML
PMLs number	12
Background index	1

a circular ring converts to an elliptic ring, LSPR wavelength is confronted to the blue shift as well as higher electric field enhancement. With  $p = 0.5$ , the nano-elliptic ring is appropriate for the waveguide to operate at  $\lambda = 1310$  nm. Notice that  $p = 1$  is represented as the circular ring and its LSPR is consistent with previous works [5].

Figure 3 demonstrates LSPR wavelength for the nano-elliptic ring with  $LD = 290$  nm,  $D = 35$  nm, and  $T = 30$  nm. Here the nano-elliptic ring with  $p = 0.5$  is selected for the waveguide to operate at  $\lambda = 1550$  nm. Hence, Figs. 2 and 3 demonstrate that the proposed plasmonic waveguide can operate at the O-band or the C-band, depending on the geometrical parameters. The FDTD simulation parameters have been listed in Table 1.

### 3. Transmission performances of gold nano-elliptic rings array in plasmonic waveguide

In this section, plasmonic waveguides are designed taking into account the Au nano-elliptic rings arrays. Electromagnetic energy transmission through the plasmonic waveguide depends on the excitation of non-collective modes ( $k \neq 0$ ) [5, 24]. Launching a temporal Gaussian pulse in time at a particular distance from the first particle, leads to the excitation of these modes. Here, the simulations are validated with waveguide

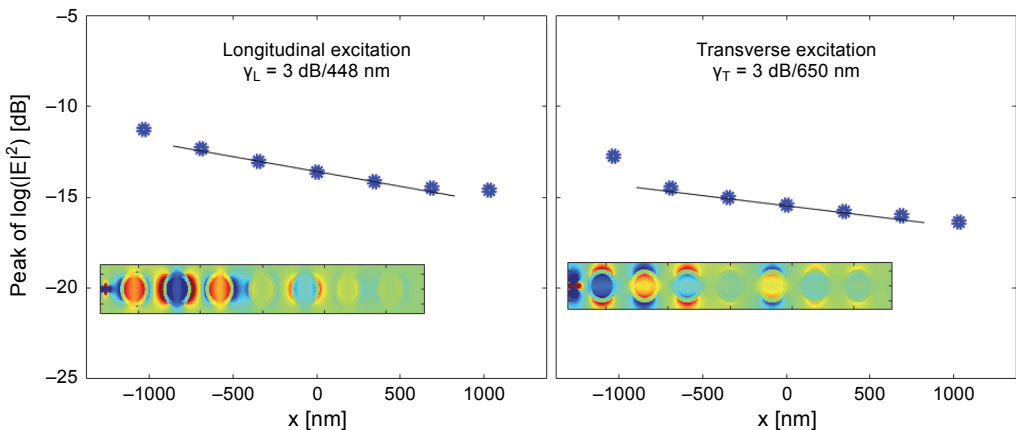


Fig. 4. Waveguide set up with the gold nanorings.

using circular nanorings arrays. Figure 4 shows that the results are consistent with previous works [5, 25]. Figure 5 demonstrates the electric field distributions in the logarithmic scale for the nano-elliptic rings, with geometrical parameters which are listed in Table 2. With the use of a linear fitting technique, except for the first and the last particle, in order to reduce the error, the waveguide transmission losses are obtained for different modes. The intensity of the propagating field decays according to an exponential function, *i.e.*,  $|E|^2 \propto \exp(-\gamma x)$ , where  $x$  is the propagation direction and  $\gamma$  is the transmission loss coefficient [5]. We obtained longitudinal and transverse  $-3$  dB trans-

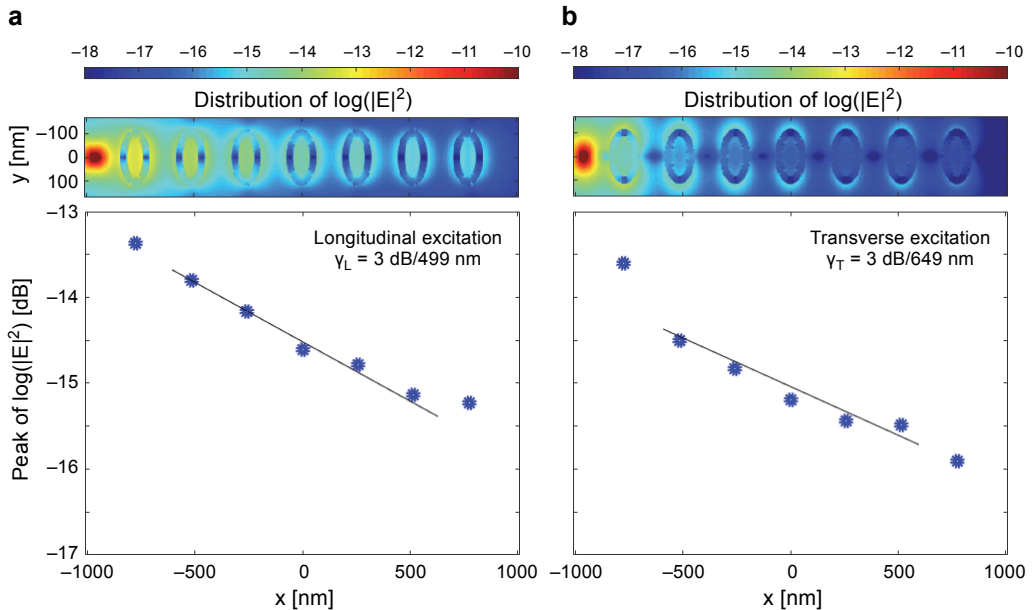


Fig. 5. Field distribution while propagating along a linear array of seven gold nano-elliptic rings chain in  $\text{SiO}_2$ , which is used to calculate the longitudinal and transverse transmission loss coefficients  $\gamma_L$  and  $\gamma_T$  at  $\lambda = 1310$  nm. The value of  $\gamma_L$  equals to 3 dB/499 nm in longitudinal excitation (a). The value of  $\gamma_T$  equals to 3 dB/649 nm in transverse excitation (b).

Table 2. Geometrical parameters for nano-elliptic arrays that operated at  $\lambda = 1310$  nm.

Description	Values [nm]
Large diameter LD	235
Small diameter SD	117.5
$p = \text{SD/LD}$	0.5
Height $H$	35
Thickness $T$	30
Inter-center distance between two particles	$1.5\text{LD} = 352.5$
Distance between source and first particle	$1.0\text{LD} = 235$

mission loss coefficients of  $\gamma_L = 3 \text{ dB}/499 \text{ nm}$  and  $\gamma_T = 3 \text{ dB}/649 \text{ nm}$  at  $\lambda = 1310 \text{ nm}$ , respectively.

Figure 6 shows the electric field distributions in the logarithmic scale for the nano-elliptic rings with geometrical parameters, which are listed in Table 3. The simulation results illustrate longitudinal and transverse  $-3 \text{ dB}$  transmission loss coefficients of  $\gamma_L = 3 \text{ dB}/737 \text{ nm}$  and  $\gamma_T = 3 \text{ dB}/791 \text{ nm}$  at  $\lambda = 1310 \text{ nm}$ , respectively. In comparison with circular rings, the elliptic rings also show minor losses, both of the modes loss coefficients are close to each other [5].

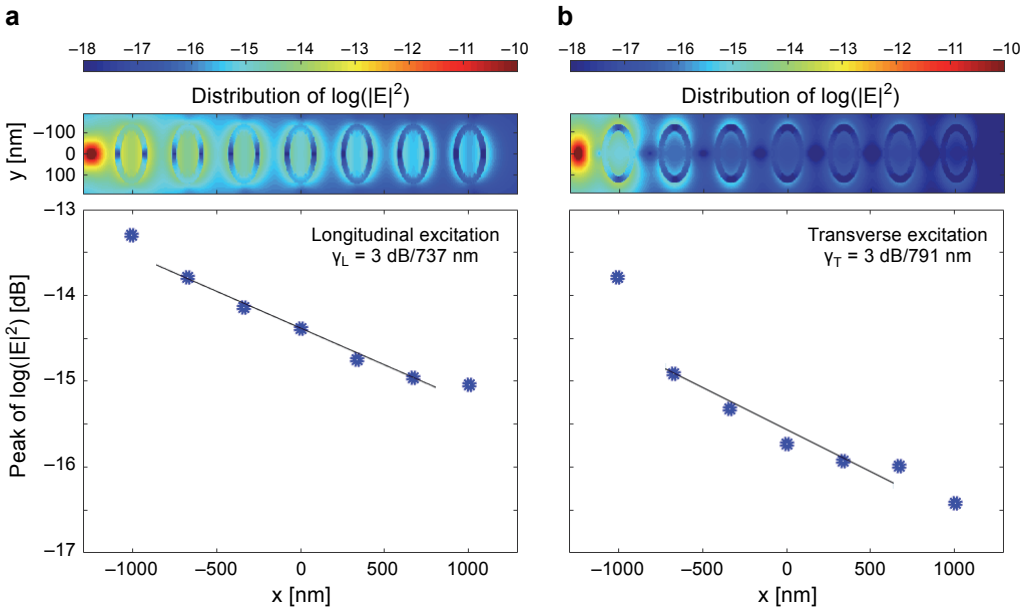


Fig. 6. Field distribution while propagating along a linear array of seven gold nano-elliptic rings chain in  $\text{SiO}_2$ , which is used to calculate the longitudinal and transverse transmission loss coefficients  $\gamma_L$  and  $\gamma_T$  at  $\lambda = 1550 \text{ nm}$ . The value of  $\gamma_L$  equals  $3 \text{ dB}/737 \text{ nm}$  in longitudinal excitation (a). The value of  $\gamma_T$  equals  $3 \text{ dB}/791 \text{ nm}$  in transverse excitation (b).

Table 3. Geometrical parameters for nano-elliptic arrays that operated at  $\lambda = 1550 \text{ nm}$ .

Description	Values [nm]
Large diameter LD	290
Small diameter SD	145
$p = \text{SD}/\text{LD}$	0.5
Height $H$	35
Thickness $T$	30
Inter-center distance between two particles	$1.5\text{LD} = 435$
Distance between source and first particle	$\text{LD} = 290$

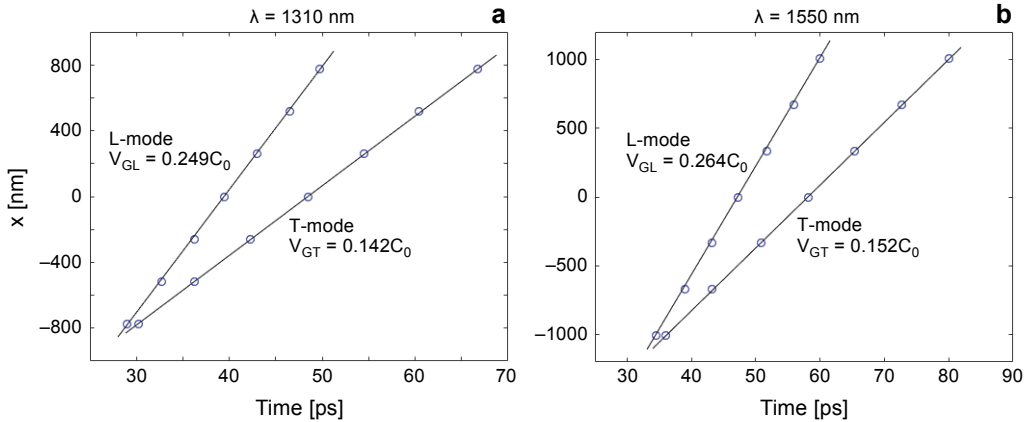


Fig. 7. Pulse peak positions over time in the plasmon waveguide consisting of arrays of nano-elliptic rings:  $V_{GL} = 0.249C_0$ , for longitudinal mode and  $V_{GT} = 0.142C_0$  for transverse mode at  $\lambda = 1310$  nm (a), and  $V_{GL} = 0.264C_0$ , for longitudinal mode and  $V_{GT} = 0.152C_0$  for transverse mode at  $\lambda = 1550$  nm (b);  $C_0$  – speed of light in vacuum.

Figure 7 shows pulse peak positions over time in the plasmon waveguide consisting of arrays of nano-elliptic rings. The longitudinal and transverse group velocities for nano-elliptic rings are a little smaller than nano-circular rings [5].

## 4. Conclusions

Three-dimensional simulations employing the finite-difference time-domain algorithm are used to study and analyze the plasmonic waveguides made up of arrays of gold nano-elliptic rings in a  $\text{SiO}_2$  host, in near infrared spectrum. Furthermore, the optical properties of an isolated gold nano-elliptic ring were studied to determine the set of geometrical parameters producing LSPR at 1310 and 1550 nm (original and communication bands). It has been found that a nano-elliptic ring with 235 nm large diameter, 117.5 nm small diameter, 30 nm metallic thickness, and 35 nm height in the  $\text{SiO}_2$  host material generates LSPR at 1310 nm wavelength and elliptic ring with 290 nm large diameter, 145 nm small diameter, 30 nm metallic thickness, and 35 nm height generates LSPR at 1550 nm wavelength with higher electric field enhancements. The transmission loss and group velocity were obtained for both longitudinal and transverse polarized excitation. At  $\lambda = 1310$  nm, the longitudinal mode was found to be more lossy than that at  $\lambda = 1550$  nm. However, at both of wavelengths, the longitudinal modes are more lossy than transverse modes, where loss coefficients in nano-elliptic ring waveguide are smaller than that in nano-circular rings waveguide, while approximately equal loss coefficients are obtained for longitudinal and transverse modes at  $\lambda = 1550$  nm. Moreover, the group velocities in the longitudinal mode are faster than the transverse

mode at both wavelengths; while nano-elliptic rings behave a little slower than circular nanorings.

## References

- [1] CHIA-YANG TSAI, JYUN-WEI LIN, CHE-YAO WU, PIN-TSO LIN, TSAN-WEN LU, PO-TSUNG LEE, *Plasmonic coupling in gold nanoring dimers: observation of coupled bonding mode*, [Nano Letters 12\(3\), 2012, pp. 1648–1654.](#)
- [2] SALEH B.E., TEICH M.C., SALEH B.E., *Fundamentals of Photonics*, Vol. 22, Wiley, New York, 1991.
- [3] RAETHER H., *Surface Plasmons on Smooth Surfaces*, Springer, 1988.
- [4] KREIBIG U., VOLLMER M., *Theoretical Considerations in Optical Properties of Metal Clusters*, Springer, 1995.
- [5] JUNG K.-Y., TEIXEIRA F.L., REANO R.M., *Au/SiO<sub>2</sub> nanoring plasmon waveguides at optical communication band*, *Journal of Lightwave Technology* **25(9)**, 2007, pp. 2757–2765.
- [6] MAIER S.A., *Plasmonics: Fundamentals and Applications*, Springer Science and Business Media, 2007.
- [7] BARNES W.L., DEREUX A., EBBESEN T.W., *Surface plasmon subwavelength optics*, [Nature 424\(6950\), 2003, pp. 824–830.](#)
- [8] OZBAY E., *Plasmonics: merging photonics and electronics at nanoscale dimensions*, [Science 311\(5758\), 2006, pp. 189–193.](#)
- [9] AIZPURUA J., BRYANT G.W., RICHTER L.J., GARCÍA DE ABAJO F.J., KELLEY B.K., MALLOUK T., *Optical properties of coupled metallic nanorods for field-enhanced spectroscopy*, [Physical Review B 71\(23\), 2005, article ID 235420.](#)
- [10] GOPINATH A., BORISKINA S.V., RANJITH PREMASIRI W., ZIEGLER L., REINHARD B.M., DAL NEGRO L., *Plasmonic nanogalaxies: multiscale aperiodic arrays for surface-enhanced Raman sensing*, [Nano Letters 9\(11\), 2009, pp. 3922–3929.](#)
- [11] THEISS J., PAVASKAR P., ECHTERNACH P.M., MULLER R.E., CRONIN S.B., *Plasmonic nanoparticle arrays with nanometer separation for high-performance SERS substrates*, [Nano Letters 10\(8\), 2010, pp. 2749–2754.](#)
- [12] TRIPATHY S., MARTY R., LIN V.K., SIEW LANG TEO, ENYI YE, ARBOUET A., SAVIOT L., GIRARD C., MING YONG HAN, MLAYAH A., *Acousto-plasmonic and surface-enhanced Raman scattering properties of coupled gold nanospheres/nanodisk trimers*, [Nano Letters 11\(2\), 2011, pp. 431–437.](#)
- [13] JUAN M.L., RIGHINI M., QUIDANT R., *Plasmon nano-optical tweezers*, [Nature Photonics 5\(6\), 2011, pp. 349–356.](#)
- [14] WURTZ G.A., POLLARD R., HENDREN W., WIEDERRECHT G.P., GOSZTOLA D.J., PODOLSKIY V.A., ZAYAT A.V., *Designed ultrafast optical nonlinearity in a plasmonic nanorod metamaterial enhanced by nonlocality*, [Nature Nanotechnology 6\(2\), 2011, pp. 107–111.](#)
- [15] LAL S., LINK S., HALAS N.J., *Nano-optics from sensing to waveguiding*, [Nature Photonics 1\(11\), 2007, pp. 641–648.](#)
- [16] SEUNG YONG LEE, LING HUNG, LANG G.S., CORNETT J.E., MAYERGOYZ I.D., RABIN O., *Dispersion in the SERS enhancement with silver nanocube dimers*, [ACS Nano 4\(10\), 2010, pp. 5763–5772.](#)
- [17] PASQUALE A.J., REINHARD B.M., NEGRO L.D., *Engineering photonic–plasmonic coupling in metal nanoparticle necklaces*, [ACS Nano 5\(8\), 2011, pp. 6578–6585.](#)
- [18] LINLIN WU, CHUNSHENG SHI, LIANGFEI TIAN, JIN ZHU, *A one-pot method to prepare gold nanoparticle chains with chitosan*, [The Journal of Physical Chemistry C 112\(2\), 2008, pp. 319–323.](#)
- [19] HAO E., SCHATZ G.C., *Electromagnetic fields around silver nanoparticles and dimers*, [The Journal of Chemical Physics 120\(1\), 2004, pp. 357–366.](#)
- [20] FENG HAO, LARSSON E.M., ALI T.A., SUTHERLAND D.S., NORDLANDER P., *Shedding light on dark plasmons in gold nanorings*, [Chemical Physics Letters 458\(4–6\), 2008, pp. 262–266.](#)



- [21] AIZPURUA J., HANARP P., SUTHERLAND D.S., KÄLL M., BRYANT G.W., GARCÍA DE ABAJO F.J., *Optical properties of gold nanorings*, [Physical Review Letters 90\(5\), 2003, article ID 057401](#).
- [22] RAYFORD C.E., SCHATZ G., SHUFORD K., *Optical properties of gold nanospheres*, *Nanoscape* 2(1), 2005, pp. 27–33.
- [23] STOLERU V., TOWE E., *Optical properties of nanometer-sized gold spheres and rods embedded in anodic alumina matrices*, [Applied Physics Letters 85\(22\), 2004, pp. 5152–5154](#).
- [24] MAIER S.A., BRONGERSMA M.L., KIK P.G., MELTZER S., REQUICHA A.A.G., KOEL B.E., ATWATER H.A., *Plasmonics—A route to nanoscale optical devices (Advanced Materials, 2001, 13, 1501)*, [Advanced Materials 15\(7–8\), 2003, pp. 562–562](#).
- [25] QUINTEN M., LEITNER A., KRENN J.R., AUSSENEGG F.R., *Electromagnetic energy transport via linear chains of silver nanoparticles*, [Optics Letters 23\(17\), 1998, pp. 1331–1333](#).

*Received September 28, 2017  
in revised form October 27, 2017*