

HEATING SOURCE LOCALIZATION IN A REDUCED TIME

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Inverse three-dimensional heat conduction problems devoted to heating source localization are ill posed. Identification can be performed using an iterative regularization method based on the conjugate gradient algorithm. Such a method is usually implemented off-line, taking into account observations (temperature measurements, for example). However, in a practical context, if the source has to be located as fast as possible (e.g., for diagnosis), the observation horizon has to be reduced. To this end, several configurations are detailed and effects of noisy observations are investigated.

Keywords: parameter identification, inverse heat conduction problem, optimal observation, source localization.

1. Introduction

In numerous engineering contexts, parameter identification is a crucial requirement when unknown parameters have to be determined in order to better understand or to control the occurring phenomena. As an illustration, let us mention, e.g., the identification of a thermal characteristic of a plasma in a nuclear fusion reactor (Mechhoud *et al.*, 2013), the identification of boundary conditions (a temperature dependent heat transfer coefficient) in a plasma assisted chemical vapor deposition process (Rouquette *et al.*, 2007a), the localization of weak heat sources in electronic devices (Rakotoniaina *et al.*, 2002) or the electron beam welding process (Rouquette *et al.*, 2007b). Among industrial applications, several usual objectives can be mentioned: predictive model validation, identification of systems dynamic behavior, estimation of thermo-physical properties, or diagnosis of processes or materials. In this paper, the proposed study is devoted to the localization of one or several stationary heating sources in a reduced time for a three-dimensional geometry. This non-linear problem can be formulated as an inverse heat conduction problem (IHCP) where the

system thermal state is described by a set of partial differential equations (PDEs). Such inverse problems are ill posed in the sense of Hadamard (Alifanov *et al.*, 1995; Isakov, 1998; Morozov, 1994). In numerous references, parameter estimation for heat flux identification is achieved considering numerical solution of inverse problems (Park *et al.*, 1999; Zhou *et al.*, 2010; Mohammadiun *et al.*, 2011; Hasanov and Pektas, 2013). Implementation of an extended Kalman filter for online identification purposes related to the IHCP was proposed by Daouas and Radhouani (2004; 2007).

Such a technique, well known in signal processing, induces a time-lag on identified parameters and can be difficult to tune without *a priori* information (related to the noise distribution, for example). Configurations studied in both previous references are related to boundary heat flux identification considering a 1D mathematical model. In the works of Kolodziej *et al.* (2010) or Mierzwiczak and Kolodziej (2010; 2011; 2012), the method of fundamental solutions was successfully implemented for an IHCP. This meshless method is relatively new and is an attractive alternative to the finite or boundary element methods. In the work of Egger *et al.* (2009), recent investigations in a 3D domain were performed in order to identify a surface heat flux using 36 microthermocouples located

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quite closely underneath the investigated surface. A conjugate gradient method is employed for iterative criterion minimization and a weak formulation is based on the Galerkin finite-dimensional technique. A multilevel approach is mentioned in order to consider sequential identification for the acceleration of minimizing the Tikhonov functional (but there are no details related to numerical implementation or tests).

In the last decades, several methods have been successfully implemented in order to determine the location of a single point heating source or line heating sources in one or two-dimension geometry. For example, in the work of Silva Neto and Ozisik (1993), both a depth and heat flux of a single heating source were identified in a one-dimension geometry using the conjugate gradient algorithm. Then, in the work of Khachfe and Jarny (2000), heat fluxes and locations of two pointwise sources in a two-dimensional domain were identified.

Le Niliot and Lefèvre (2001) investigated identification of multiple steady line heat sources in a diffusive system and proposed its application in a 2D experimental configuration. Always focused on 2D geometry, Lefèvre and Le Niliot (2002) estimated the location and strength of line heat sources using an identification procedure based on a boundary integral formulation with transient fundamental solutions. In the work of Yi and Murio (2002), a source term identification in a 1D-IHCP is performed using a regularization procedure, based on the mollification method and a marching scheme for the solution of the stabilized problem.

Based on an earlier reference (Beck and Arnold, 1977) and using the boundary element method, parameter identification was successfully investigated (Le Niliot and Lefèvre, 2004) in order to solve an inverse problem of pointwise heat source identification. Recently, a two-step regularized inverse solution for two-dimensional heat source reconstruction was developed by Renault *et al.* (2008). Moreover, the authors proposed a spectral method for the estimation of a thermo-mechanical heat source (this method is based on the analytical solution of the direct problem in the Laplace–Fourier domain and applied to a two-dimensional problem of heat conduction) (Renault *et al.*, 2010). Girault *et al.* (2010) implemented the modal identification method for the identification of both the location and the heat flux of a heating source in a two-dimensional heat diffusion problem.

In the following, the method proposed by Silva Neto and Ozisik (1993) as well as Khachfe and Jarny (2000) is extended to a three-dimensional geometry. The IHCP is solved in order to determine the locations of one or several stationary heating sources while assuming that heat fluxes are known. It is important to notice that the spatial distribution of heat sources is not pointwise and has to be taken into account in a small area of

the domain surface. Source locations are determined given temperature evolutions observed by a few pointwise sensors. In this context and among various methods available to solve the IHCP, the iterative conjugate gradient method (CGM) has been implemented for the minimization of the quadratic cost function describing the error between the numerical model outputs and the sensor measurements (Hager and Zhang, 2006; Prudhomme and Nguyen, 1998). It is well known that for such iterative minimization the solution of three well posed problems (always in the sense of Hadamard) is required. The direct problem describing evolution of the system state (temperature) is solved in order to evaluate the quadratic cost function. The gradient of the cost function is obtained once the adjoint problem is solved. Then the sensitivity problem is considered for the determination of the step-length in a descent direction. As for noisy measurements, it is shown by Alifanov (1994) that the conjugate gradient method acts as a regularization method. In fact, a stopping criterion (for the minimization algorithm) which takes into account noise standard deviation can be considered a regularization parameter. In the proposed study, localization in a reduced time is investigated in order to put in evidence the relation between noise levels, the number of unknown parameters and the duration of measurements.

This paper is organized as follows. In the next section, the investigated thermal system is exposed (the notation and the PDE system). In Section 3, the IHCP is formulated and a brief description of the conjugate gradient algorithm is proposed. In order to locate the heating sources, both the sensitivity and adjoint problems are presented. In Section 4, several numerical examples are provided for different situations: one, two, and three sources. A reduced time solution is investigated and the effect of noisy measurements is presented. Finally, conclusion and outlooks are briefly drawn up at the end of this paper.

2. Direct problem

This section is devoted to the description of thermal heating phenomena modeled by a set of partial differential equations. Let us consider a titanium square plate $\Omega \subset \mathbb{R}^3$ (length $L = 5 \cdot 10^{-2}$ m, thickness $e = 10^{-3}$ m); the domain boundary is denoted by $\partial\Omega \subset \mathbb{R}^2$ (see Fig. 1). The space variable is

$$(x, y, z) \in \left] -\frac{L}{2}, \frac{L}{2} \right[\times \left] -\frac{e}{2}, \frac{e}{2} \right[\times \left] -\frac{L}{2}, \frac{L}{2} \right[$$

and the time variable is $t \in T = [0; t_f]$ in seconds. The plate is heated on its lower face $\Gamma_{\text{heated}} = \{(x, -e/2, z) \in \partial\Omega\}$ by one or several heating sources S_j ($j = 1, \dots, N_s$, where N_s is the number of heating sources). For each source, the heat flux is denoted by ϕ_j

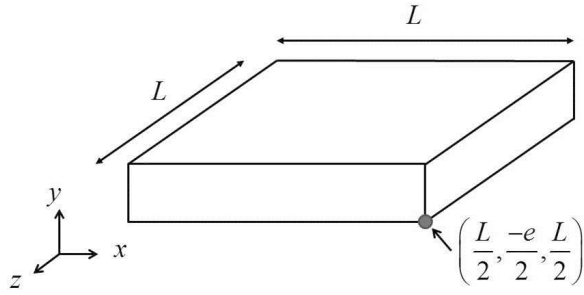


Fig. 1. Studied plate geometry.

and is not time dependent. Moreover, the heat flux ϕ_j is assumed to be uniform on the disk $D_j \in \Gamma_{\text{heated}}$. Each disk (source) center is denoted by $I_j = (x_j, -e/2, z_j)$ and the disk radius is $r = 0.002$ m. From a practical point of view, these circular heating flux spatial distributions could be obtained using a proper optical homogenizer. Such an optical device (also called the kaleidoscope) is commonly used to obtain a uniform flux (Chen *et al.*, 1963) on a heated surface. Museux *et al.* (2012) used such equipment in order to obtain a uniform flux on a square for evaluation of burns. However, specific kaleidoscopes can be easily designed in order to obtain a circular distribution of a radiative heating flux on a plane target. Let us consider the following model for heating flux density in $W \cdot m^{-2}$:

$$\Phi(x, z) = \sum_{j=1}^{N_s} \phi_j F_j(x, z),$$

where

$$F_j(\cdot) = \frac{1}{\pi} \left(-\text{atan}(\mu(\sigma_j(x, z) - r)) + \frac{\pi}{2} \right),$$

with $\mu \in \mathbb{R}^+$ and

$$\sigma_j(x, z) = \sqrt{(x - x_j)^2 + (z - z_j)^2}.$$

The parameter μ is chosen in order to describe the discontinuity in the neighborhood of the disk boundary. For the proposed configuration, $\mu = 10^5$ is reasonable (image analysis is easily performed). It is important to notice that the parametrized function leads to a description of heat fluxes in a continuous and differentiable way.

Moreover, a natural convection phenomenon is considered on all plate boundaries $\partial\Omega$: the convection coefficient is h in $W \cdot m^{-2} \cdot K^{-1}$. The 3D geometry considered is presented in Fig. 1. A more complex heat transfer and semi-transparent media are examined for predictive purposes in (Lormel *et al.*, 2004). The temperature evolution is described by the following equations:

Direct problem P_{dir} : Assuming a uniform initial temperature θ_0 (equal to the ambient

temperature) and the set of known parameters $P = \{\lambda, \rho c, h, \Phi, L, e\}$, determine the temperature $\theta(x, y, z; t)$ as the solution of

$$\left\{ \begin{array}{l} \forall (x, y, z; t) \in \Omega \times T, \\ \lambda \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} + \frac{\partial^2 \theta}{\partial z^2} \right) = \rho c \frac{\partial \theta}{\partial t}, \\ \forall (x, y, z) \in \Omega, \quad \theta(x, y, z; 0) = \theta_0, \\ \forall (x, y, z; t) \in \partial\Omega \times T, \\ -\lambda \left(\frac{\partial \theta}{\partial \vec{n}} \right) = h(\theta - \theta_0) - \Phi, \end{array} \right. \quad (1)$$

where \vec{n} is the outward normal unit vector.

Model input parameters are presented in Table 1. If all model inputs parameters are fixed, the previous direct problem described by (1) can be numerically solved using the finite element method (implemented, for example, in the Comsol-Multiphysics software interfaced with Matlab). In the following, the temperature $\theta(x, y, z; t)$ numerically predicted at given spatial locations will be considered as the measured temperature for a set of sensors (see Section 3). In the proposed setting, the sensor spatial location is chosen as follows. It is crucial to select enough measurements points. It is obvious that a single sensor will be unable to find more than a single source (if heating fluxes are constant). The proposed configuration with three sensors is enough to allow triangulation. Since the heating source locations are unknown, in order to investigate the sample surface without a favored area, an equilateral triangle centered on the upper plate is meaningful. This is not a systematic procedure based on theoretical hypothesis about the investigated system a but experimental reasoning. If a large number of sensors is available, an optimal selection can be implemented based on the methodology developed by Perez and Vergnaud (2016). In the next section, the conjugate gradient method (CGM) is implemented to solve inverse problems in order to identify source locations I_j in a reduced time.

3. Inverse problem

3.1. Formulation of the inverse heat conduction problem (IHCP).

Assume that the coordinates of the source

Table 1. Notation and parameter definitions.

Square length	$L = 5 \cdot 10^{-2}$ m
Plate thickness	$e = 2 \cdot 10^{-3}$ m
Initial temperature	$\theta_0 = 293$ K
Heat flux	$\phi_j = 10^5$ $W \cdot m^{-2}$
Thermal conductivity	$\lambda = 21.9$ $W \cdot m^{-1} \cdot K^{-1}$
Volume heat	$\rho c = 2.35 \cdot 10^6$ $J \cdot m^{-3} \cdot K^{-1}$
Natural convection coefficient	$h = 20$ $W \cdot m^{-2} \cdot K^{-1}$

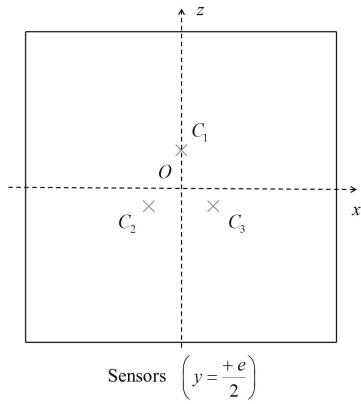


Fig. 2. Sensor location.

centers

$$I_j^* \left(x_j^*, \frac{-e}{2}, z_j^* \right)_{j=1, \dots, N_s}$$

are unknown. In the proposed setting, the aim of IHCP solution is to identify the source coordinates (x_j^*, z_j^*) on the lower face given temperatures obtained from three sensors located on the opposite face $y = +e/2$. This IHCP can be formulated as a classic optimization problem where a quadratic criterion denoted by $J(I)$ has to be minimized. Let us denote by $\hat{\theta}(t)$ the temperature measured at three sensors located on the upper face:

$$\begin{aligned} C_1 & \left(0, \frac{+e}{2}, 0.006 \right), \\ C_2 & \left(-0.0052, \frac{+e}{2}, -0.003 \right), \\ C_3 & \left(0.0052, \frac{+e}{2}, -0.003 \right) \end{aligned}$$

(see Fig. 2). In the framework of optimal sensor location, several thermal situations encountered for high temperature processes are presented by Autrique *et al.* (2000; 2002). The quadratic criterion $J(I)$ is formulated as

$$J(I) = \frac{1}{2} \int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I) - \hat{\theta}_i(t) \right)^2 dt, \quad (2)$$

where $\theta(C_i; t; I)$ is the temperature predicted for sensor i (given the source location I) while $\hat{\theta}_i(t)$ is the temperature measured at sensor i . The cost function is discretized using sampled measurements with a time step Δt (N_t measurements are available for each sensor):

$$J_d(I) = \frac{\Delta t}{2} \sum_{n=1}^{N_t} \sum_{i=1}^3 \left(\theta(C_i; t_n; I) - \hat{\theta}_i(t_n) \right)^2.$$

The inverse problem is as follows:

Inverse problem P_{inv} : Find unknown source locations $I_{j=1, \dots, N_s}^*$ such that the quadratic functional $J(I)$ is minimum:

$$I^* = \arg \min_{\mathbb{R}^{2N_s}} J(I), \quad (3)$$

subject to the constraint $\theta(x, y, z; t; I)$ being a solution of the direct problem P_{dir} (1).

In order to solve the previous ill-posed problem, the CGM algorithm is implemented (see, e.g., Alifanov, 1994; Silva Neto and Ozisik, 1993; Tarantola, 2005). It is important to note that the descent direction is reset to $(\beta^k = 0)$ every $2N_s$ iterations. Such a restart procedure for conjugate directions is discussed by Powell (1977). The direct problem (1) is discussed in Section 2. The following sections are concerned with the presentation of the sensitivity and adjoint problems related to Steps 3 and 4 of the CGM algorithm (see Algorithm 1).

3.2. Sensitivity problem. The formulation of this problem is helpful for the statement of the adjoint problem and for the estimation of the step-length in the descent direction. Let $\delta\theta(x, y, z; t)$ be the temperature variation induced by a variation in the unknown parameters $\delta I = \{\delta x_j, \delta y_j\}_{j=1, \dots, N_s}$. Applications of variational calculus are presented by Weinstock (1952). For the PDE satisfied by the temperature $\theta(x, y, z; t) + \epsilon\delta\theta(x, y, z; t)$ (see the direct problem (1) with source locations $I + \epsilon\delta I$), as $\epsilon \rightarrow 0$, the sensitivity problem (Khachfe and Jarny, 2000; 2001) becomes

Sensitivity problem P_{sens} : Given the set of parameters $P = \{\lambda, \rho c, h, \delta\Phi, L, e\}$, determine $\delta\theta(x, y, z; t)$ as the solution of

$$\begin{cases} \forall(x, y, z; t) \in \Omega \times T, \\ \lambda \left(\frac{\partial^2 \delta\theta}{\partial x^2} + \frac{\partial^2 \delta\theta}{\partial y^2} + \frac{\partial^2 \delta\theta}{\partial z^2} \right) = \rho c \frac{\partial \delta\theta}{\partial t}, \\ \forall(x, y, z) \in \Omega, \quad \delta\theta(x, y, z; 0) = 0, \\ \forall(x, y, z; t) \in \partial\Omega \times T, \\ -\lambda \left(\frac{\partial \delta\theta}{\partial n} \right) = h\delta\theta - \delta\Phi. \end{cases} \quad (4)$$

Here we have

$$\delta\Phi(x, z) = \sum_{j=1}^{N_s} \phi_j \left(\mathbf{1}_{D_j^+ \setminus D_j} - \mathbf{1}_{D_j \setminus D_j^+} \right)$$

where

$$\mathbf{1}_{A \setminus B}(x, z) = \begin{cases} 1 & \text{if } (x, z) \in A \text{ and } (x, z) \notin B, \\ 0 & \text{otherwise,} \end{cases}$$

and D_j^+ is the disk with radius r and heating source center $I_j + \epsilon \delta I_j$. This heating flux variation can be also formulated considering

$$\delta \Phi(x, z) = \sum_{j=1}^{N_s} \phi_j \delta F_j(x, z)$$

with

$$\begin{aligned} \delta F_j(x, z) &= \frac{1}{\pi} \frac{-\mu(\delta x_j(x_j - x) + \delta z_j(z_j - z))}{\sigma_j(x, z) \left(1 + (\mu(\sigma_j(x, z) - r))^2\right)} \\ &= A_j(x, z) \delta x_j + B_j(x, z) \delta z_j, \end{aligned}$$

where

$$\begin{aligned} A_j(x, z) &= \frac{1}{\pi} \frac{-\mu(x_j - x)}{\sigma_j(x, z) \left(1 + (\mu(\sigma_j(x, z) - r))^2\right)}, \\ B_j(x, z) &= \frac{1}{\pi} \frac{-\mu(z_j - z)}{\sigma_j(x, z) \left(1 + (\mu(\sigma_j(x, z) - r))^2\right)}. \end{aligned}$$

Finally, the heating flux variation is expressed as

$$\delta \Phi(x, z) = \sum_{j=1}^{N_s} \phi_j (A_j(\cdot) \delta x_j + B_j(\cdot) \delta z_j). \quad (5)$$

3.3. Adjoint problem. This problem consists in determining the gradient of the cost function

$$\overrightarrow{\nabla J^k} = \left(\frac{\partial J}{\partial x_j}, \frac{\partial J}{\partial z_j} \right)_{j=1, \dots, N_s}$$

at each iteration k of the CGM by introducing an adjoint function $\psi(x, y, z; t)$. Other applications are presented by Jarny *et al.* (1991), Silva Neto and Ozisik (1993), Huang and Chen (1999), Khachfe and Jarny (2000; 2001), Rouquette *et al.* (2007a) and Perez *et al.* (2007).

Let us introduce the Lagrangian $\mathcal{L}(\theta, I, \psi)$ associated with the direct problem (1):

$$\begin{aligned} \mathcal{L}(\theta, I, \psi) &= J(\theta(\cdot)) \\ &+ \int_0^{t_f} \int_{\Omega} \left(\rho c \frac{\partial \theta(\cdot)}{\partial t} - \lambda \Delta \theta(\cdot) \right) \psi(\cdot) dt d\Omega. \end{aligned} \quad (6)$$

Let us consider

$$\delta \mathcal{L}(\theta, I, \psi) = \frac{\partial \mathcal{L}}{\partial \theta} \delta \theta + \frac{\partial \mathcal{L}}{\partial I} \delta I + \frac{\partial \mathcal{L}}{\partial \psi} \delta \psi$$

with

$$\frac{\partial \mathcal{L}}{\partial I} \delta I = \sum_{j=1}^{N_s} \left(\frac{\partial \mathcal{L}}{\partial x_j} \delta x_j + \frac{\partial \mathcal{L}}{\partial z_j} \delta z_j \right).$$

Algorithm 1. CGM algorithm.

Step 1. Initialization $k = 0$: $I_j^0 (x_j^0, z_j^0)_{j=1, \dots, N_s}$ is the initial unknown source center.

Step 2. Direct problem solution in order to calculate the temperature and the cost function $J(I^k)$.

- If $J(I^k) \leq J_{\text{stop}}$ (admissible threshold J_{stop} will be discussed further), the minimization procedure is stopped and the current values $I_j^k (x_j^k, z_j^k)_{j=1, \dots, N_s}$ are considered adequate estimates for source coordinates I_j^* ;
- otherwise, the iterative procedure is continued.

Step 3. Adjoint problem solution in order to compute the functional gradient

$$\overrightarrow{\nabla J^k} = \left(\frac{\partial J}{\partial x_j}, \frac{\partial J}{\partial z_j} \right) (x_j^k, z_j^k)$$

for $j = 1, \dots, N_s$, and the next descent direction

$$\overrightarrow{d^{k+1}} = -\overrightarrow{\nabla J^k} + \beta^k \overrightarrow{d^k},$$

where

$$\beta^k = \frac{\|\overrightarrow{\nabla J^k}\|^2}{\|\overrightarrow{\nabla J^{k-1}}\|^2}$$

($\|\cdot\|$ is the Euclidean norm and $\beta^0 = 0$).

Step 4. Sensitivity problem solution in order to estimate the step-length in the descent direction:

$$\gamma^{k+1} = \arg \min_{\gamma \in \mathbb{R}} (I^k - \gamma \overrightarrow{d^{k+1}})$$

Step 5. Estimation of the new coordinates $I_j^{k+1} (x_j^{k+1}, z_j^{k+1})$ using the following formulas:

$$x_j^{k+1} = x_j^k - \gamma^{k+1} \left(\overrightarrow{d_j^{k+1}} \right)_1$$

and

$$z_j^{k+1} = z_j^k - \gamma^{k+1} \left(\overrightarrow{d_j^{k+1}} \right)_2,$$

where

$$\left(\left(\overrightarrow{d_j^{k+1}} \right)_1, \left(\overrightarrow{d_j^{k+1}} \right)_2 \right)$$

corresponds to the descent direction related to heat source j and coordinates (x, z) .

Step 6. $k \leftarrow k + 1$ and go to Step 2.

If θ is a solution to (1), then $\mathcal{L}(\theta, I, \psi) = J(I)$ and $\delta\mathcal{L}(\theta, I, \psi) = \delta J(\theta, I)$. If ψ is fixed, then

$$\frac{\partial \mathcal{L}}{\partial \psi} \delta \psi = 0$$

and $\delta\mathcal{L}(\theta, I, \psi)$ becomes

$$\delta\mathcal{L}(\theta, I, \psi) = \frac{\partial \mathcal{L}}{\partial \theta} \delta \theta + \sum_{j=1}^{N_s} \left(\frac{\partial \mathcal{L}}{\partial x_j} \delta x_j + \frac{\partial \mathcal{L}}{\partial z_j} \delta z_j \right).$$

Then the choice of the fixed Lagrange multiplier $\psi(x, y, z; t)$ is performed in order to satisfy the following equation:

$$\frac{\partial \mathcal{L}}{\partial \theta} \delta \theta = 0. \tag{7}$$

In order to determine $\psi(x, y, z; t)$ which satisfies (7), while taking account of the previous equation (6), let us formulate $\delta\mathcal{L}(\theta, I, \psi)$ as follows:

$$\begin{aligned} \delta\mathcal{L}(\theta, I, \psi) &= \int_0^{t_f} \int_{\Omega} E \delta \theta \, dt \, d\Omega \\ &+ \int_0^{t_f} \int_{\Omega} \left(\rho c \frac{\partial \delta \theta}{\partial t} - \lambda \Delta \delta \theta \right) \psi \, dt \, d\Omega, \end{aligned}$$

where

$$\begin{aligned} E(x, y, z; t) &= \sum_{i=1}^3 \left(\theta(x, y, z; t) - \hat{\theta}_i(t) \right) \delta_d(x, y, z; C_i) \end{aligned}$$

and $\delta_d(x, y, z; C_i)$ is the Dirac distribution which is considered at the sensor location C_i . After several integrations by parts (versus space and time), while taking account of the sensitivity equations (4), Eqn. (7) is satisfied if the Lagrange multiplier ψ is a solution to the following problem:

Adjoint problem P_{adj} : Assuming the set of parameters $P = \{\lambda, \rho c, h, L, e, E\}$, determine $\psi(x, y, z; t)$ as a solution to

$$\left\{ \begin{aligned} &\forall (x, y, z; t) \in \Omega \times T, \\ &\lambda \left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) = E - \rho c \frac{\partial \psi}{\partial t}, \\ &\forall (x, y, z) \in \Omega, \quad \psi(x, y, z; t_f) = 0, \\ &\forall (x, y, z; t) \in \partial\Omega \times T, \quad -\lambda \left(\frac{\partial \psi}{\partial \vec{n}} \right) = h\psi. \end{aligned} \right. \tag{8}$$

The corresponding derivation is presented in Appendix. If ψ is a solution to the adjoint problem P_{adj} (8), then

$$\delta\mathcal{L}(\theta, I, \psi) = - \int_0^{t_f} \int_{\partial\Omega} \psi \delta\Phi \, dt \, d(\partial\Omega),$$

with $\delta\Phi$ given in (5). Then $\delta\mathcal{L}(\theta, I, \psi)$ can be written as

$$\begin{aligned} \delta\mathcal{L}(\theta, I, \psi) &= - \int_0^{t_f} \int_{\partial\Omega} \psi \sum_{j=1}^{N_s} \frac{\phi_j(t)}{\pi} (A_j(\cdot) \delta x_j \\ &+ B_j(\cdot) \delta z_j) \, dt \, d(\partial\Omega). \end{aligned}$$

Since

$$\langle \vec{\nabla} J, \delta I \rangle_{L^2(T, \partial\Omega)} = \delta J = \delta\mathcal{L},$$

the gradient of the cost function is

$$\begin{aligned} \vec{\nabla} J &= \left(\frac{\partial J}{\partial x_j}, \frac{\partial J}{\partial z_j} \right)_{j=1, \dots, N_s} \\ &= \left(- \int_0^{t_f} \int_{\partial\Omega} \frac{\phi_j}{\pi} A_j \psi \, dt \, d(\partial\Omega), \right. \\ &\quad \left. - \int_0^{t_f} \int_{\partial\Omega} \frac{\phi_j}{\pi} B_j \psi \, dt \, d(\partial\Omega) \right)_{j=1, \dots, N_s}. \end{aligned} \tag{9}$$

3.4. Step-length in the descent direction. The step-length in the descent direction γ^{k+1} corresponds to the optimal depth in the descent direction for the new value of the unknown heating source coordinates:

$$\begin{aligned} I^{k+1} &= I^k - \gamma^{k+1} \overrightarrow{d^{k+1}} \\ \gamma^{k+1} &= \arg \min_{\mathbb{R}} J(I^{k+1}), \end{aligned} \tag{10}$$

where

$$\begin{aligned} J(I^{k+1}) &= \frac{1}{2} \int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I^k - \gamma^{k+1} \overrightarrow{d^{k+1}}) - \hat{\theta}_i(t) \right)^2 \, dt. \end{aligned}$$

For small variations in heating source localizations, the problem can be linearized and the temperature is written as

$$\begin{aligned} \theta(C_i; t; I^k - \gamma^{k+1} \overrightarrow{d^{k+1}}) &= \theta(C_i; t; I^k) - \gamma^{k+1} \delta\theta_{\overrightarrow{d^{k+1}}}(C_i; t). \end{aligned}$$

Since $\delta\theta_{\overrightarrow{d^{k+1}}}(C_i; t)$ is the sensitivity function induced by the variation in the heating source coordinates denoted by $\overrightarrow{d^{k+1}}$, we have

$$\begin{aligned} J(I^k - \gamma^{k+1} \overrightarrow{d^{k+1}}) &= \frac{1}{2} \int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I^k) - \hat{\theta}_i(t) \right)^2 \, dt \\ &+ \frac{(\gamma^{k+1})^2}{2} \int_0^{t_f} \sum_{i=1}^3 \left(\delta\theta_{\overrightarrow{d^{k+1}}}(C_i; t) \right)^2 \, dt \end{aligned}$$

$$\begin{aligned}
 & -\gamma^{k+1} \int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I^k) - \hat{\theta}_i(t) \right) \\
 & \times \delta \theta_{\vec{d}^{k+1}}(C_i; t) dt.
 \end{aligned}$$

In order to minimize the quadratic functional $J(I^k - \gamma^{k+1} \vec{d}^{k+1})$, γ^{k+1} is determined such that

$$\frac{\partial J(I^k - \gamma^{k+1} \vec{d}^{k+1})}{\partial \gamma} = 0,$$

and then

$$\begin{aligned}
 & \frac{\partial J(I^k - \gamma^{k+1} \vec{d}^{k+1})}{\partial \gamma} \\
 & = \gamma^{k+1} \int_0^{t_f} \sum_{i=1}^3 \left(\delta \theta_{\vec{d}^{k+1}}(C_i; t) \right)^2 dt \\
 & - \int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I^k) - \hat{\theta}_i(t) \right) \delta \theta_{\vec{d}^{k+1}}(C_i; t) dt.
 \end{aligned}$$

Then the step-length in the descent direction γ^{k+1} is

$$\begin{aligned}
 & \gamma^{k+1} \\
 & = \frac{\int_0^{t_f} \sum_{i=1}^3 \left(\theta(C_i; t; I^k) - \hat{\theta}_i(t) \right) \delta \theta_{\vec{d}^{k+1}}(\cdot) dt}{\int_0^{t_f} \sum_{i=1}^3 \left(\delta \theta_{\vec{d}^{k+1}}(\cdot) \right)^2 dt}.
 \end{aligned} \quad (11)$$

At each iteration k , the sensitivity problem P_{sens} (4) has to be solved in the descent direction \vec{d}^{k+1} to compute the step-length in the descent direction γ^{k+1} .

3.5. Admissible level of minimization. In order to implement the algorithm of the CGM, a threshold J_{stop} has to be defined in order to stop the iterative minimization procedure. If all perturbations can be neglected and the experiment takes place in ideal conditions, then, without noisy measurements and model errors, J_{stop} can be ideally fixed close to zero. In this numerical study, observations are obtained from a given direct problem (where heat source locations

$$I_j^* \left(x_j^*, \frac{-e}{2}, z_j^* \right)_{j=1, \dots, N_s}$$

are assumed to be known). These simulated observations are taken into account to determine the source locations in a reduced time. If temperature measurements are not arbitrarily disturbed, the minimization algorithm is stopped when $J(I^k, \theta) < J_{\text{stop}} = 10^{-3} \text{ K}^2 \cdot \text{s}$. For such a value, source locations are well estimated since the additional criterion (tracking error)

$$\text{TE}(I^k) = \sum_{j=1}^{N_s} \sqrt{(x_j^k - x_j^*)^2 + (z_j^k - z_j^*)^2}$$

is quite small, where $I_j^k(x_j^k, -e/2, z_j^k)$ is the estimated location of the heat source $I_j^*(x_j^*, -e/2, z_j^*)$ at iteration k . In the numerical setting considered, where I_j^* are known, TE is relevant information to investigate the algorithm behavior. The tracking error is considered a stopping criterion and the algorithm is stopped if $\text{TE}(I^k) < 10^{-3} \text{ m}$. It is important to notice that, for realistic experimental situations when heat source locations are unknown, the tracking error cannot be calculated. Then the stopping criterion can be considered equal to

$$J_{\text{stop}} = \Delta t \frac{N_c}{2} \sigma^2 N_t,$$

where σ is the standard deviation of the Gaussian noise and N_t is the number of observations for each of $N_c = 3$ sensors. Such a stopping strategy is discussed by Alifanov (1994).

4. Numerical results

In the following paragraphs, several configurations are investigated in order to illustrate CGM implementation and its regularization effect in the context of identification of heat source locations. For one, two or three heat sources, the minimization of the cost function and the evolution of the tracking error are shown for observations obtained at three sensors (with or without noise). Numerical results are obtained using the Comsol-Multiphysics solver interfaced with Matlab software with/without noisy measurements, for different values of the final time t_f (this key parameter can also be considered the necessary measurement time).

4.1. Case A.1. In this first case, the plate is heated with a single source located at $I(0.01, -e/2, 0.01)$ in meters (see Fig. 3). Assume that $t_f = 1 \text{ s}$. Numerical results for the direct problem are presented in Fig. 4. Observations are performed every 0.1 s and the temperature values numerically obtained are considered as measurements $\theta(C_i, t)$ for the set of sensors C_i (see Fig. 2). In the following, the source coordinates are unknown. In order to identify these unknown coordinates, the IHCP described by P_{inv} (see Section 3) is solved. The initial value of source localization is given by $I(0.01, -e/2, 0)$. The solution of P_{inv} is performed using CGM and results are presented hereafter. The algorithm is stopped for $J(I^k, \theta) < J_{\text{stop}} = 10^{-3}$. In Fig. 5, it is shown that after eight iterations, criterion $J(I^k, \theta)$ converges under the adopted stopping criterion. Moreover, evolution of the tracking error, which decreases at each iteration, confirms the correct behavior of the minimization algorithm. For $J(I^k, \theta) < J_{\text{stop}} = 10^{-3}$, the tracking error is lower than 10^{-3} m . Coordinate evolution (versus iteration k) of the sought source is shown in Fig. 6. The identified

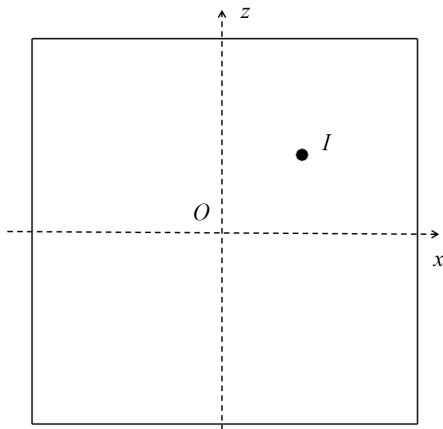


Fig. 3. Source location $y = -e/2$: Case A.1.

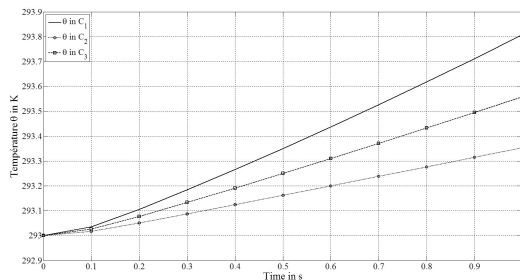


Fig. 4. Direct problem solution: Case A.1.

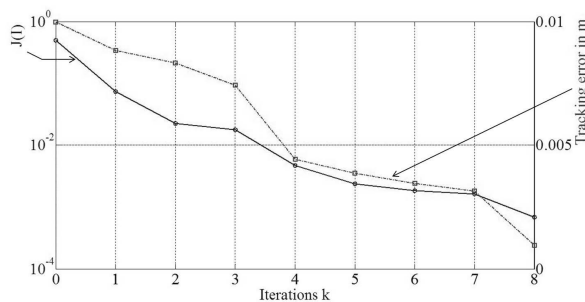


Fig. 5. Cost function and the tracking error: Case A.1.

Table 2. Criterion values vs. iteration k : Case A.1.

Iteration k	0	1	2	3	...
$J(I)$	0.50	0.07	0.023	0.018	...
Iteration k	...	7	8		
$J(I)$...	0.0016	0.0007		

coordinates versus iteration k are presented in Table 3. The error for the localization of the disk center is lower than 0.95 mm. These previous results show that heat source coordinates are obtained after eight iterations with the measurement time equal to one second. The choice of this final time for identification purposes plays a crucial role in the convergence of the conjugate

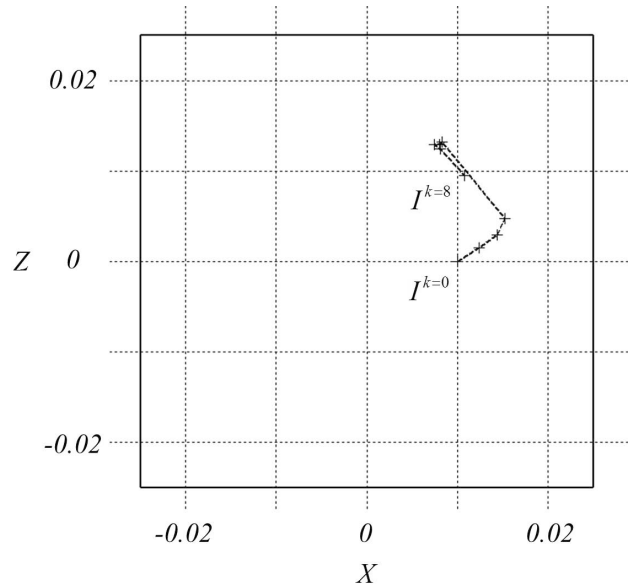


Fig. 6. Coordinate evolution on the heated face: Case A.1.

Table 3. Source coordinates (in mm) vs. iteration k : Case A.1.

Iteration k	0	1	2	3	...
x	10	12.4	14.4	15.3	...
z	0	1.5	2.9	4.8	...
Iteration k	...	7	8		
x	...	8.1	10.8		
z	...	12.5	9.5		

Table 4. Criterion and tracking error (in mm) values versus iteration k : Case A.2.

k	0	1	2	3	4
J	4.28	1.26	0.97	0.90	0.85
TE	10	9	8	5	7
k	5	6	7	...	12
J	0.75	0.729	0.727	...	0.727
TE	6	7	6.8	...	0.5

gradient algorithm. This choice depends on the number of information pieces provided by sensors. For noisy measurements, the final time has to be adapted in order to obtain the desired convergence (and the correct location).

4.2. Case A.2. The temperature predicted in the same configuration as for Case A.1 is disturbed by additive Gaussian noise defined by $\mathcal{N}(0, 0.5)$.

Both the cost-function and tracking error versus iteration k are shown in Table 4 and Fig. 7. It is shown that the convergence of the iterative algorithm is achieved in 12 iterations with the convergence of the tracking error under the threshold value $TE(I^k, \theta) < 10^{-3}$. In order to illustrate that the final time $t_f = 1s$ does not provide enough observations, tracking errors are compared for

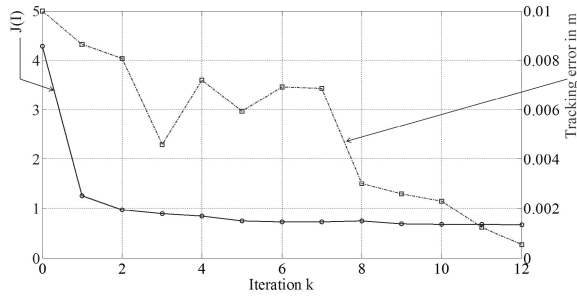


Fig. 7. Cost function and the tracking error: Case A.2.

Table 5. Source coordinates (in mm) vs. iteration k : Case A.2.

k	0	1	2	3	4
x	10	12.8	14.6	14.2	6.7
z	0	1.8	3.3	8.2	16.4
k	5	...	12		
x	5.2	...	8.9		
z	13.5	...	10.5		

$t_f = 1$ s and $t_f = 2$ s in Fig. 8. It is shown that for $t_f = 1$ s, the tracking error is almost constant for $k \geq 6$ and it is not possible to identify the correct heat source location. For $t_f = 2$ s the tracking error obtained after 12 iterations is lower than the desired accuracy. Coordinate evolution is shown in Table 5. The error for the localization of the disk center is lower than 1.21 mm.

4.3. Case A.3. Let us consider a greater standard deviation, i.e., Gaussian noise $\mathcal{N}(0, 1)$ (see Fig. 9). Both the cost-function and the tracking error versus iteration k are shown in Table 6 and Fig. 10. The identified source coordinates are presented in Table 7. The error for the localization of the disk center is lower than 0.77 mm. Moreover, in Fig. 11, evolution of the tracking error for several final time values is presented. The obtained results show that three seconds ($t_f = 3$ s) are sufficient to identify the heating source coordinates considering the measured temperature disturbed by Gaussian noise $\mathcal{N}(0, 1)$, while 2 seconds do not provide enough disturbed observations to

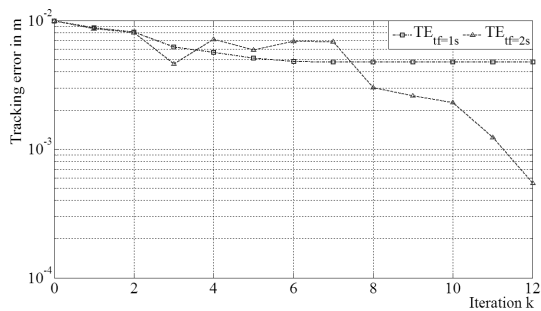


Fig. 8. Tracking error for different final time values: Case A.2.

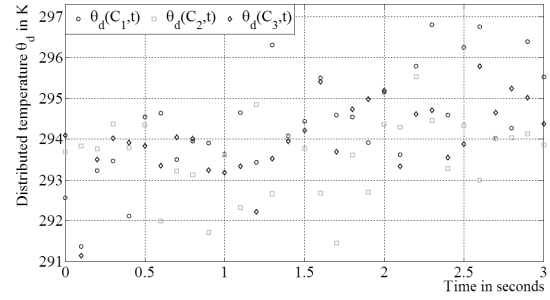


Fig. 9. Disturbed measurements for $t_f = 3$ s: Case A.3.

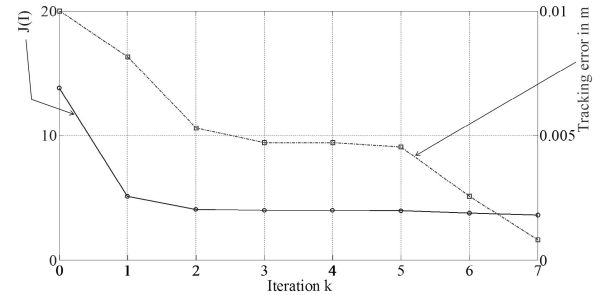


Fig. 10. Cost function and the tracking error: Case A.3.

identify the correct location.

The previous cases confirm robustness of the CGM when identifying a single heat source location. Moreover, the regularization effect of the CGM is illustrated when noisy data are taken into account for identification purposes.

Table 6. Criterion and tracking error (in mm) values vs. iteration k : Case A.3.

k	0	1	2	3	4
J	13.81	5.105	4.074	3.968	3.968
TE	10	8	5	4.7	4.7
k	5	6	7		
J	3.958	3.755	3.596		
TE	4.5	2.6	0.8		

Table 7. Source coordinates vs. iteration k : Case A.3.

k	0	1	2	3	4
x	10	12.9	13.7	12.4	12.4
z	0	2.4	6.2	5.9	5.9
k	5	6	7		
x	12	10.3	10.3		
z	6	7.4	9.3		

4.4. Case B.1. Consider two heat sources whose real positions are defined by $I_1 (0.01, -e/2, 0.01)$ and

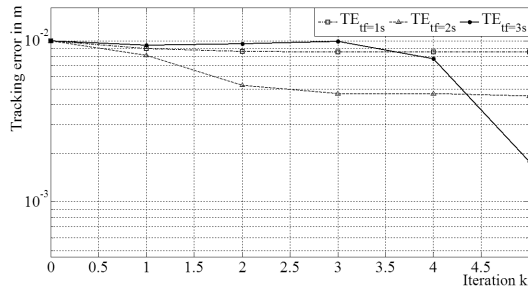


Fig. 11. Tracking error for different final time values: Case A.3.

Table 8. Criterion values vs. iteration k : Case B.1.

Iteration k	0	1	2	3	4
$J(I_1, I_2)$	1.14	0.18	0.074	0.058	0.047
Iteration k	...	8	9		
$J(I_1, I_2)$...	0.0016	$4 \cdot 10^{-5}$		

$I_2 (-0.01, -e/2, 0.01)$ in meters (see Fig. 12). Simulated temperatures (treated as measured temperatures for identification purposes) are presented in Fig. 13.

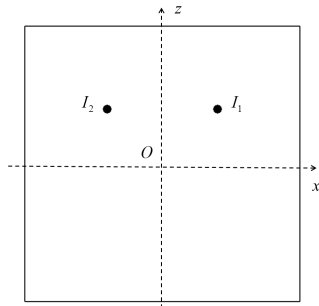


Fig. 12. Source location $y = -e/2$: Case B.1.

Assume that for the inverse problem solution the CGM is implemented with initial sources coordinates equal to $I_1^{k=0} (0.01, -e/2, 0)$ and $I_2^{k=0} (-0.01, -e/2, 0)$. Results are presented in Fig. 14 and Table 8. Evolution of the identified coordinates for the two heat sources is shown in Fig. 15 and the values are presented in Table 9. The error for the localization of the disk center is lower than 0.1 mm. In this case, observations performed during only one second ($t_f = 1$ s) are relevant to identify the locations of two sources. The effect of noisy measurements is investigated in the following section.

4.5. Case B.2. Assume that the measured temperatures are disturbed by Gaussian noise $\mathcal{N}(0, 0.5)$ (see Fig. 16). Evolution of both the cost function and the tracking error is outlined in Fig. 17 and Table 10. The obtained coordinates of the localizations of the two heating sources are given in Table 11. The error for the localization of the disk center is lower than 0.81 mm. In order to justify the

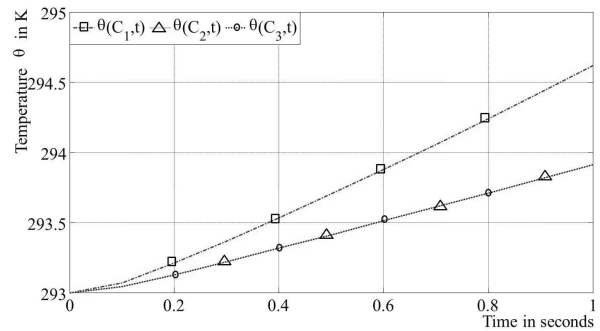


Fig. 13. Direct problem solution: Case B.1.

choice of the final time ($t_f = 2$ s), a comparison between the tracking errors obtained for $t_f = 1$ s and $t_f = 2$ s versus iteration k is presented in Fig. 17. It is shown that, for noise $\mathcal{N}(0, 0.5)$, observations performed during 1 second are not efficient for identifying the locations of two heat sources. In fact, in such a situation the signal/noise ratio is too small and the CGM fails to extract relevant information so quickly. For a final time equal to 2 seconds, convergence is correct and unknown coordinates are accurately identified.

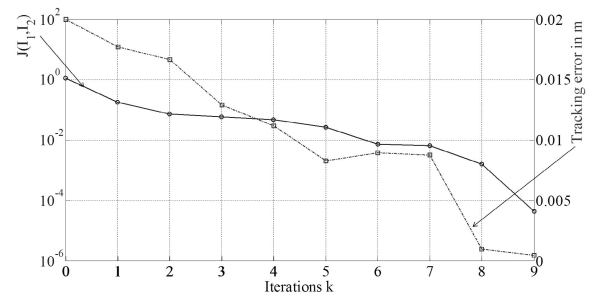


Fig. 14. Cost function and the tracking error: Case B.1.

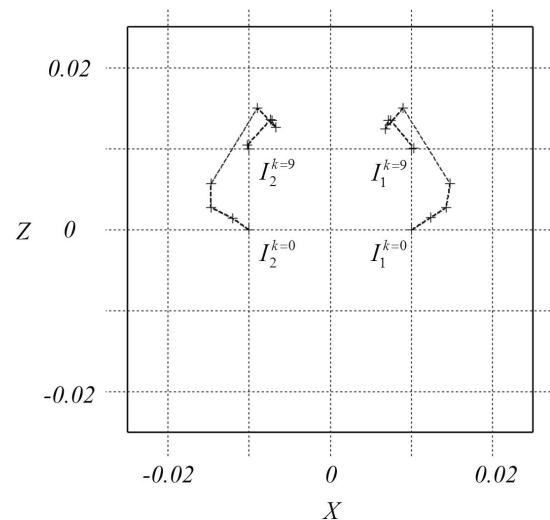


Fig. 15. Source coordinates evolution on the heated face: Case B.1.

Table 9. Source coordinates (in mm) vs. iteration k : Case B.1.

k	0	1	2	3	4
x_1	10	12.4	14.3	14.8	9
z_1	0	1.5	2.8	5.7	15
x_2	-10	-12	-14.7	-14.7	-9
z_2	0	1.4	2.8	5.7	15
k	5	6	7	8	9
x_1	6.8	7.2	7.5	10.3	10
z_1	12.5	13.5	13.5	10.4	10
x_2	-6.7	-7.2	-7.4	-10.2	-10
z_2	12.6	13.5	13.6	10.5	10

4.6. Case B.3. For a larger noise standard deviation $\mathcal{N}(0, 1)$, the predicted temperature is shown in Fig. 19. The CGM for of the identification source coordinates is called for $t_f = 10$ s and results are presented in Fig. 20 and Tables 12 and 13. The localization error of the disk center is lower than 1.21 mm. According to the

Table 11. Source coordinates (in mm) vs. iteration k : Case B.2.

k	0	...	3	4	5
x_1	10	...	12.7	7	4.1
z_1	0	...	11.5	19	13.3
x_2	-10	...	-12.3	-6	-3.3
z_2	0	...	11.5	19	12.6
k	...	10	11	12	...
x_1	...	15.3	3.5	6	...
z_1	...	8.4	17.6	15	...
x_2	...	-17.3	-1.6	-5	...
z_2	...	9.4	16.9	15	...
k	15	16	17	...	23
x_1	12	8	7.5	...	9.6
z_1	9	14.3	13.5	...	10.7
x_2	-14	-7.3	-6.7	...	-9.6
z_2	10	13.9	13.1	...	10.3

previous results with the final time $t_f = 10$ s, it is shown that the CGM is robust and efficient when identifying the

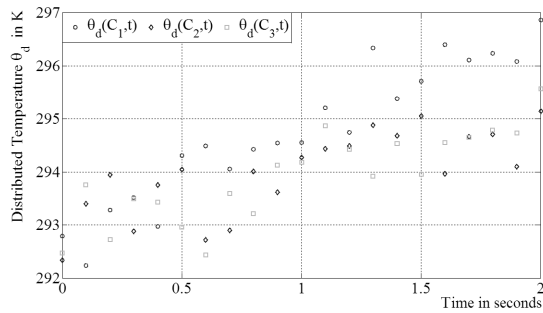


Fig. 16. Disturbed measurements for $t_f = 2$ s: Case B.2.

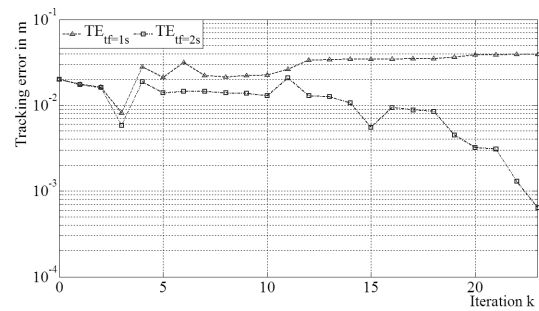


Fig. 18. Tracking error for different final times: Case B.2.

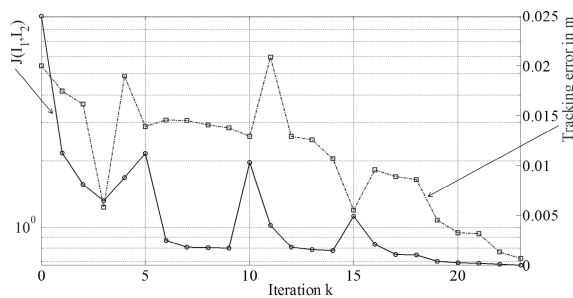


Fig. 17. Cost function and the tracking error vs. k : Case B.2.

Table 10. Cost function and tracking error (in mm) vs. k : Case B.2.

k	0	1	2	3	4
J	9.2	2.2	1.6	1.3	1.7
TE	20	17	16	6	20
k	5	6	...	23	
J	2.17	0.087	...	0.67	
TE	14	14	...	0.6	

Table 12. Cost function and the tracking error (in MM) vs. k : Case B.3.

k	0	...	3	4	5
J	384	...	50	40	27
TE	20	...	7	14	12
k	...	7	8	...	11
J	...	26.37	26.37	...	12.84
TE	...	8.5	3.5	...	0.5

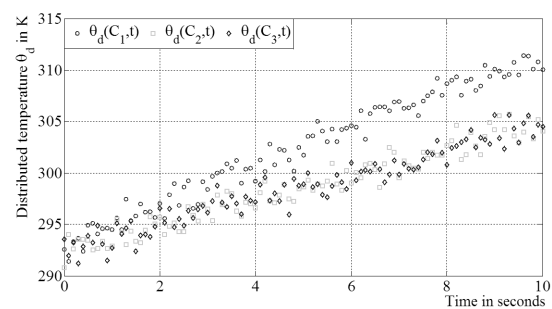


Fig. 19. Disturbed measurements for $t_f = 10$ s: Case B.3.

locations of two heat sources. In order to justify the choice $t_f = 10$ s, several final times are compared in Fig. 21. It is important to note that for observations performed during 10 seconds, even with significant noise (see Fig. 19), heat source coordinates are accurately identified in only 11 iterations.

4.7. Case C.1. For this last studied configuration, the CGM is implemented in order to identify the location of three heat sources (S_1 , S_2 and S_3) located at $I_1 (0.01, -e/2, 0.01)$, $I_2 (-0.01, -e/2, 0.01)$ and $I_3 (0, -e/2, -0.005)$ in meters (see Fig. 22).

Simulated temperatures (treated as the measured temperatures for identification purposes) are presented in Fig. 23. The corresponding initial source coordinates are $I_1^{k=0} (0.01, -e/2, 0)$, $I_2^{k=0} (-0.01, -e/2, 0)$ and $I_3^{k=0} (0, -e/2, 0)$.

The results are presented in Figs. 24 and 25 and Tables 14 and 15. The error for the localization of the disk center is lower than 0.23 mm. For this situation, if the final time is lower than 7 s, then the algorithm cannot converge and the cost function cannot decrease under the stopping criterion fixed as $J_{stop} = 10^{-3}$ (see Section 3.5). In fact, an important amount of observations is required in order to discriminate between the three sources.

4.8. Case C.2. Measurements disturbed by Gaussian noise $\mathcal{N}(0, 0.5)$ are considered in Fig. 26. The results are shown in Fig. 27 and Tables 16 and 17. The error for the localization of the disk center is lower than 0.65 mm.

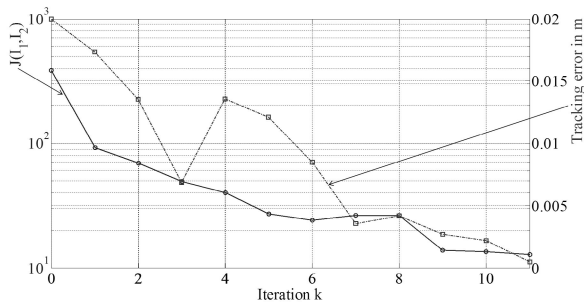


Fig. 20. Cost function and tracking error evolution: Case B.3.

Table 13. Source coordinates (in mm) vs. iteration k : Case B.3.

k	0	1	2	3	4
x_1	10	13.8	15	11	6.4
z_1	0	2.2	5	13	15.7
x_2	-10	-13.8	-15	-10	-6.1
z_2	0	2.2	5	13	15.6
k	5	6	...	11	
x_1	5.8	7.9	...	9.5	
z_1	14.3	13.8	...	11.1	
x_2	-5.66	-7.8	...	-9.5	
z_2	14.2	13.6	...	10.8	

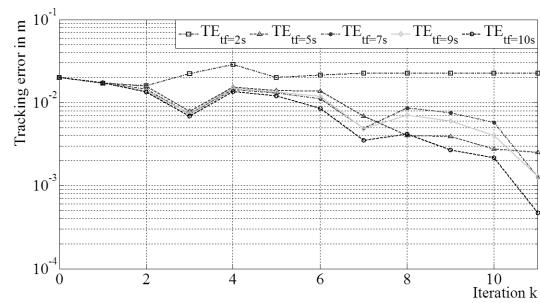


Fig. 21. Tracking error for different final times: Case B.3.

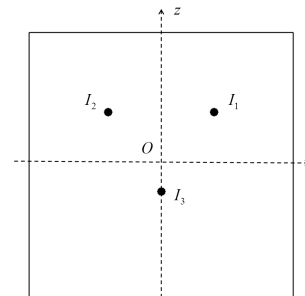


Fig. 22. Source location $y = -e/2$: Case C.1.

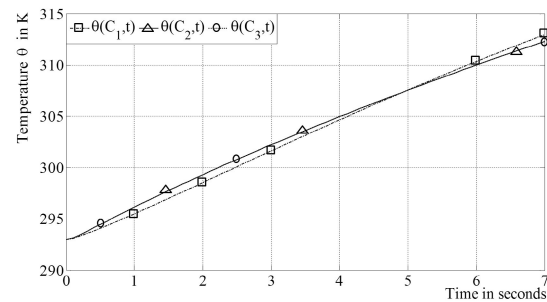


Fig. 23. Direct problem solution: Case C.1.

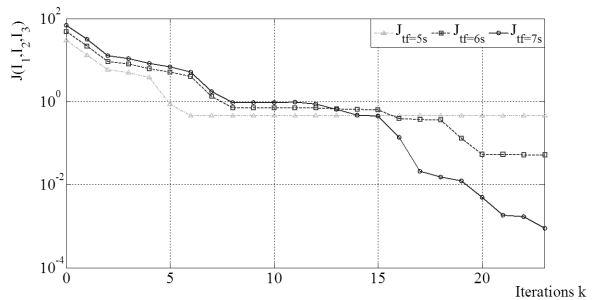


Fig. 24. Cost function and the tracking error vs. k : Case C.1.

Table 14. Criterion values vs. iteration k : Case C.1.

k	0	1	2	...	22	23
J	70	32	13	...	0.002	0.0009

In Fig. 28, evolution of the tracking error is shown for several final times, and it is evidenced that a time longer than for non-disturbed observations is required for

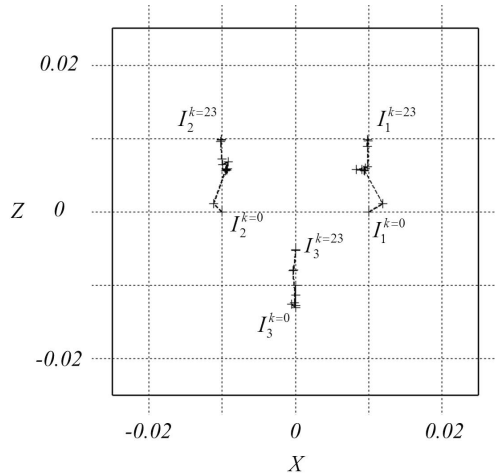


Fig. 25. Evolution of the source coordinates on the heated face: Case C.1.

Table 15. Source coordinates (in mm) vs. iteration k : Case C.1.

k	0	1	2	...	17	...	23
x_1	10	11.9	9.5	...	9.9	...	9.8
z_1	0	1.1	5.6	...	9.7	...	9.9
x_2	-10	-11.2	-9.6	...	-10.2	...	-10.1
z_2	0	1.1	5.6	...	9.6	...	9.9
x_3	0	0	-0.01	...	0	...	0
z_3	-10	-11.3	-12.4	...	-5.1	...	-5.2

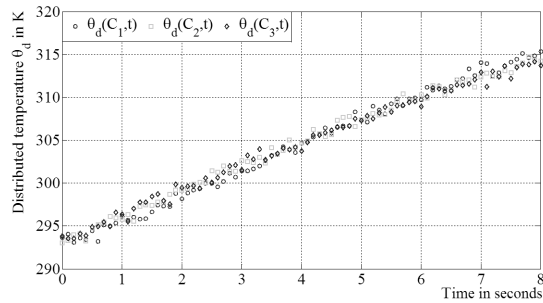


Fig. 26. Disturbed measurements for $t_f = 8$ s: Case C.2.

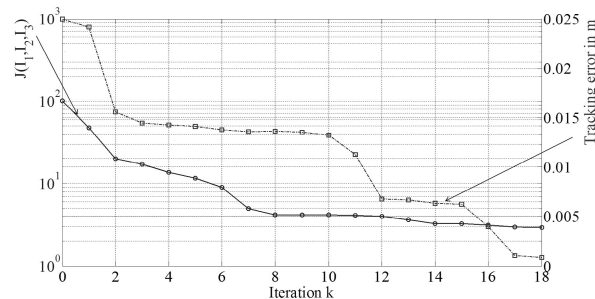


Fig. 27. Cost function and the tracking error: Case C.2.

location identification.

4.9. Case C.3. For this final case, additive Gaussian noise $\mathcal{N}(0, 1)$ is considered, and three sensors and three heat sources have to be located. The measurements are

Table 16. Cost function and the tracking error (in mm) vs. iteration k : Case C.2.

k	0	1	2	3	4	...	18
J	101	47	20	17	14	...	2.96
TE	25	24	16	14.4	14.2	...	0.8

Table 17. Source coordinates (in mm) versus iteration k : Case C.2.

k	0	1	2	...	11	12
x_1	10	11.3	9.5	...	9.6	10
z_1	0	1.2	6.0	...	6.3	7.6
x_2	-10	-11.1	-9.3	...	-9.2	-9.6
z_2	0	1.2	5.9	...	6.2	7.4
x_3	0	0	0	...	0	-0.07
z_3	-10	-10	-11.4	...	-8.8	-6.8

k	...	15	...	17	18
x_1	...	9.6	...	9.9	9.9
z_1	...	8.1	...	8.9	10.2
x_2	...	-9.9	...	-10.9	-10.5
z_2	...	7.9	...	9.4	9.6
x_3	...	0.09	...	-0.01	0.02
z_3	...	-7.2	...	-5.3	-5.1

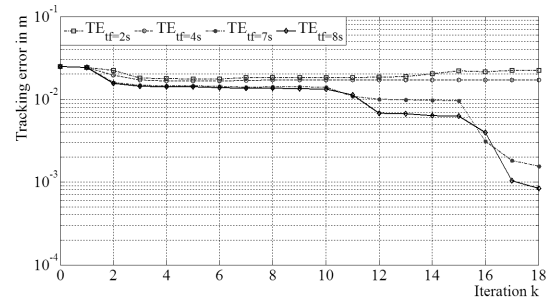


Fig. 28. Tracking error evolution for different final time values: Case C.2.

shown in Fig. 29. The results are shown in Fig. 30 as well as Tables 18 and 19 for the final time $t_f = 16$ s. The error for the localization of the disk center is lower than 0.45 mm. In Fig. 31, evolution of the tracking error is shown for several final times.

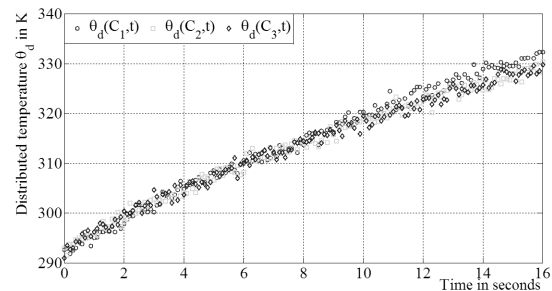


Fig. 29. Disturbed measurements for $t_f = 16$ s: Case C.3.

If $t_f < 16$ s, the tracking error does not converge under the desired threshold.

5. Analysis

According to the previous results, the necessary measurement time t_f depends on the standard deviation of the noise and $2N_s$ unknown parameters. This can be illustrated by Table 20. In realistic cases, where the tracking error TE cannot be estimated, it is more difficult to define the reduced time t_f . It seems to be attractive to investigate the signal/noise ratio and to compare the energy signal

$$\mathcal{E} = \int_0^{t_f} \sum_{i=1}^3 (\hat{\theta}_i(t) - \theta_0)^2 dt$$

with the admissible level of minimization

$$J_{stop} = \frac{3}{2} \sigma^2 n_t.$$

It is important to notice that both of these functions are increasing in time and are obtained without any *a priori* information on the source location. Without noise, it is important to obtain enough information in order to discriminate between sources. The final time depends on the distance between sources and sensors, as well as the heat flux. In the three studied configurations, this time can be described by $t_f = 1 + t_0$, where t_0 is defined as the time for which $\mathcal{E} > e^{7.5(N_s-2)}$. If noise has to be taken into account, then $t_f = 1 + t_0 + t_1$, where t_1 is defined as the time for which $\mathcal{E} > (J_{stop})^{0.6N_s}$.

It is obvious that the previous formulas are heuristically determined and adapted to the studied configurations. However, the complexity problem is taken into account: t_0 depends on the number of sources and t_1 is related to the number of observations and the standard

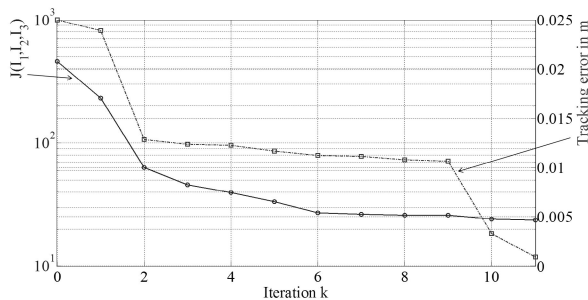


Fig. 30. Cost function and the tracking error: Case C.3.

Table 18. Cost function and the tracking error (in mm) vs. iteration k : Case C.3.

k	0	1	2	3	4	...	11
J	457	231	63	46	40	...	23.6
TE	25	24	13	12.3	12.2	...	0.9

Table 19. Source coordinates (in mm) vs. iteration k : Case C.3.

k	0	1	2	...	9	10	11
x_1	10	11.5	8.9	...	9.0	9.4	9.6
z_1	0	1.4	7.2	...	7.0	9.1	9.9
x_2	-10	-11.5	-9.0	...	-9.2	-9.1	-10
z_2	0	1.5	7.1	...	7	9.1	9.9
x_3	0	0	0.01	...	0.02	0.4	-0.4
z_3	-10	-11.5	-11.8	...	-9.4	-6.3	-5.2

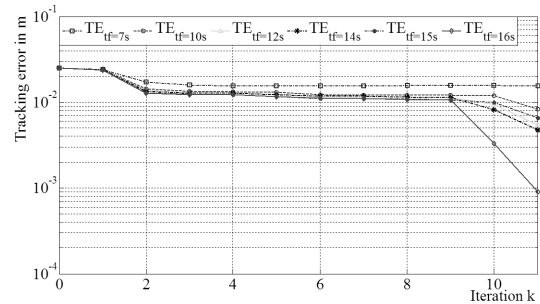


Fig. 31. Tracking error for different final times t_f : Case C.3.

Table 20. Required measurement time t_f (in s).

	$\sigma = 0$	$\sigma = 0.5$	$\sigma = 1$
$N_s = 1$	1	2	3
$N_s = 2$	1	2	10
$N_s = 3$	7	8	16

Table 21. Admissible level of minimization and signal energy (in $K^2 \cdot s$).

J_{stop}	$\sigma = 0$	$\sigma = 0.5$	$\sigma = 1$
$N_s = 1$	0	7.5	45
$N_s = 2$	0	7.5	150
$N_s = 3$	0	30	240
\mathcal{E}	$\sigma = 0$	$\sigma = 0.5$	$\sigma = 1$
$N_s = 1$	0.4	3	10
$N_s = 2$	1.4	11	1400
$N_s = 3$	2800	4200	33400

Table 22. Predicted t_f (in s).

	$\sigma = 0$	$\sigma = 0.5$	$\sigma = 1$
$N_s = 1$	1	3	4
$N_s = 2$	2	4	7
$N_s = 3$	7	8	17

deviation of the noise. In Table 22, predicted values $t_f = 1 + t_0 + t_1$ are given and can be compared with Table 20.

6. Conclusion

In this paper, an inverse heat conduction problem (IHCP) is formulated in order to estimate the location of one or several heat sources in a three-dimensional

geometry using the conjugate gradient method (CGM). For this algorithm implementation, the solution of three well-posed problems is required. Numerical experiments have been performed in order to study the effect of the number of sources, measurement noise, regularization, and the observation time required for accurate identification. Robustness of the proposed algorithm has been illustrated. If an adequate number of observations is available, identification is achieved in several iterations. This adequate number of observations is directly connected to the final time. In experimental situations, it has to be determined considering the signal/noise ratio, the number of sources and the admissible level of minimization.

Several directions can be considered for further works. Theoretical determination of the required observation time has to be investigated in specific academic situations. Simultaneous identification of both the location and the strength of fixed heat sources is also an interesting goal considering the necessary measurement time. Moreover, for mobile heat source tracking, a sliding time window based on minimal time determination seems to be quite an attractive approach. Last but not least, an experimental device will be developed in the LARIS-Lab research institute in order to test numerous situations for several heat sources.

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Appendix

In order to better understand the formulation of the boundary and final conditions in the adjoint problem, consider the variation in the Lagrangian

$$\begin{aligned} \delta\mathcal{L}(\theta, I, \psi) &= \int_0^{t_f} \int_{\Omega} E\delta\theta \, dt \, d\Omega \\ &+ \int_0^{t_f} \int_{\Omega} \left(\rho c \frac{\partial\delta\theta}{\partial t} - \lambda\Delta\delta\theta \right) \psi \, dt \, d\Omega \\ &= \mathcal{I}_1 + \mathcal{I}_2. \end{aligned}$$

The second term is developed using integration by parts in time and space:

$$\mathcal{I}_2 = \int_0^{t_f} \int_{\Omega} \left(\rho c \frac{\partial\delta\theta}{\partial t} - \lambda\Delta\delta\theta \right) \psi \, dt \, d\Omega$$

$$\begin{aligned} &= \int_0^{t_f} \int_{\Omega} \rho c \frac{\partial\delta\theta}{\partial t} \psi \, dt \, d\Omega - \int_0^{t_f} \int_{\Omega} \lambda\Delta\delta\theta \psi \, dt \, d\Omega, \\ &= \mathcal{I}_3 - \mathcal{I}_4, \\ \mathcal{I}_3 &= \int_0^{t_f} \int_{\Omega} \rho c \frac{\partial\delta\theta}{\partial t} \psi \, dt \, d\Omega \\ &= \int_{\Omega} \left([\rho c \delta\theta \psi]_0^{t_f} - \int_0^{t_f} \rho c \delta\theta \frac{\partial\psi}{\partial t} \, dt \right) d\Omega. \end{aligned}$$

Setting the initial condition $\delta\theta(x, y, z; 0) = 0$ in the sensitivity problem, we have

$$\mathcal{I}_3 = \int_{\Omega} \rho c \delta\theta \psi(\cdot; t_f) \, d\Omega - \int_0^{t_f} \int_{\Omega} \rho c \delta\theta \frac{\partial\psi}{\partial t} \, dt \, d\Omega$$

For the last term, we get

$$\begin{aligned} \mathcal{I}_4 &= \int_0^{t_f} \int_{\Omega} \lambda\Delta\delta\theta \psi \, dt \, d\Omega \\ &= \int_0^{t_f} \left(\int_{\partial\Omega} \lambda \frac{\partial\delta\theta}{\partial \vec{n}} \psi \, d\delta\Omega - \int_{\partial\Omega} \lambda \delta\theta \frac{\partial\psi}{\partial \vec{n}} \, d\delta\Omega \right) dt \\ &+ \int_0^{t_f} \int_{\Omega} \lambda \delta\theta \Delta\psi \, d\Omega \, dt \end{aligned}$$

Setting the boundary condition for the sensitivity problem

$$\forall (x, y, z; t) \in \partial\Omega \times T, \quad -\lambda \left(\frac{\partial\delta\theta}{\partial \vec{n}} \right) = h\delta\theta - \delta\Phi,$$

we have

$$\begin{aligned} \mathcal{I}_4 &= \int_0^{t_f} \int_{\partial\Omega} (-h\delta\theta + \delta\Phi) \psi \, d\delta\Omega \, dt \\ &- \int_0^{t_f} \int_{\partial\Omega} \lambda \delta\theta \frac{\partial\psi}{\partial \vec{n}} \, d\delta\Omega \, dt \\ &+ \int_0^{t_f} \int_{\Omega} \lambda \delta\theta \Delta\psi \, d\Omega \, dt \\ &= - \int_0^{t_f} \int_{\partial\Omega} \left(h\psi + \lambda \frac{\partial\psi}{\partial \vec{n}} \right) \delta\theta \, d\delta\Omega \, dt \\ &+ \int_0^{t_f} \int_{\partial\Omega} \delta\Phi \psi \, d\delta\Omega \, dt + \int_0^{t_f} \int_{\Omega} \lambda \delta\theta \Delta\psi \, d\Omega \, dt. \end{aligned}$$

Consequently, the variation in the Lagrangian becomes

$$\begin{aligned} \delta\mathcal{L} &= \int_0^{t_f} \int_{\Omega} E\delta\theta \, dt \, d\Omega \\ &+ \int_{\Omega} \rho c \delta\theta \psi(\cdot; t_f) \, d\Omega - \int_0^{t_f} \int_{\Omega} \rho c \delta\theta \frac{\partial\psi}{\partial t} \, dt \, d\Omega \end{aligned}$$

$$\begin{aligned}
 & + \int_0^{t_f} \int_{\partial\Omega} \left(h\psi + \lambda \frac{\partial\psi}{\partial\vec{n}} \right) \delta\theta \delta\Omega \, dt \\
 & - \int_0^{t_f} \int_{\partial\Omega} \delta\Phi\psi \, d\delta\Omega \, dt - \int_0^{t_f} \int_{\Omega} \lambda\delta\theta\Delta\psi \, d\Omega \, dt \\
 = & \int_0^{t_f} \int_{\Omega} \left(E - \rho c \frac{\partial\psi}{\partial t} - \lambda\Delta\psi \right) \delta\theta \, dt \, d\Omega \\
 & + \int_{\Omega} \rho c \delta\theta\psi(\cdot; t_f) \, d\Omega - \int_0^{t_f} \int_{\partial\Omega} \delta\Phi\psi \, d\delta\Omega \, dt \\
 & + \int_0^{t_f} \int_{\partial\Omega} \left(h\psi + \lambda \frac{\partial\psi}{\partial\vec{n}} \right) \delta\theta \delta\Omega \, dt.
 \end{aligned}$$

As for Eqn. (7), the Lagrange multiplier ψ has to be a solution of the following system of PDEs:

$$\begin{cases} \forall(x, y, z; t) \in \Omega \times T, \\ \lambda \left(\frac{\partial^2\psi}{\partial x^2} + \frac{\partial^2\psi}{\partial y^2} + \frac{\partial^2\psi}{\partial z^2} \right) = E - \rho c \frac{\partial\psi}{\partial t}, \\ \forall(x, y, z) \in \Omega, \quad \psi(x, y, z; t_f) = 0, \\ \forall(x, y, z; t) \in \partial\Omega \times T, \quad -\lambda \left(\frac{\partial\psi}{\partial\vec{n}} \right) = h\psi. \end{cases}$$

Thus,

$$\delta\mathcal{L}(\theta, I, \psi) = - \int_0^{t_f} \int_{\partial\Omega} \psi \delta\Phi \, dt \, d(\partial\Omega).$$

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